



Modeling for spectroscopic diagnostics of CCP-ICP and PBIF plasmas in KAERI

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2023. 7. 12



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- Collisional-radiative modeling (CRM) for He, H/D plasmas
 - Modeling methods and AM data
 - Sensitivities to used AM data and plasma parameters
- S/XB ratio for sputtering yield of W
 - S/XB ratio methods and atomic data
 - Experimental S/XB ratios

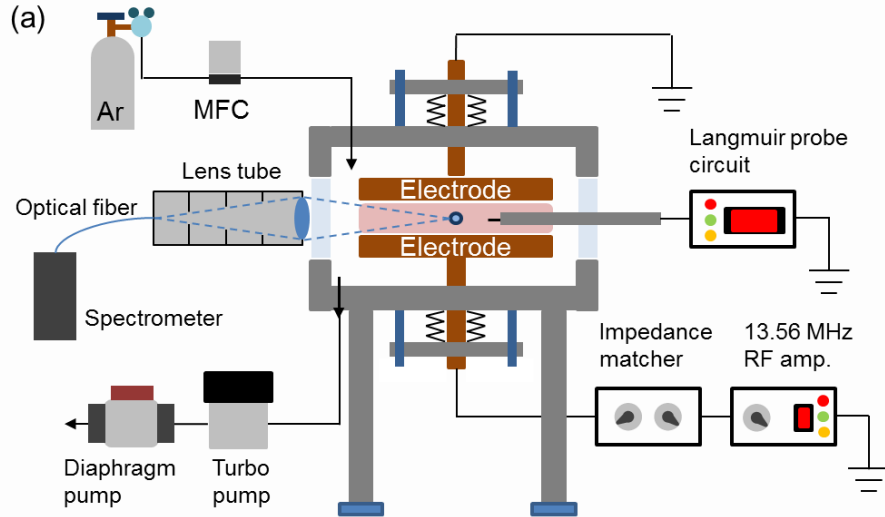
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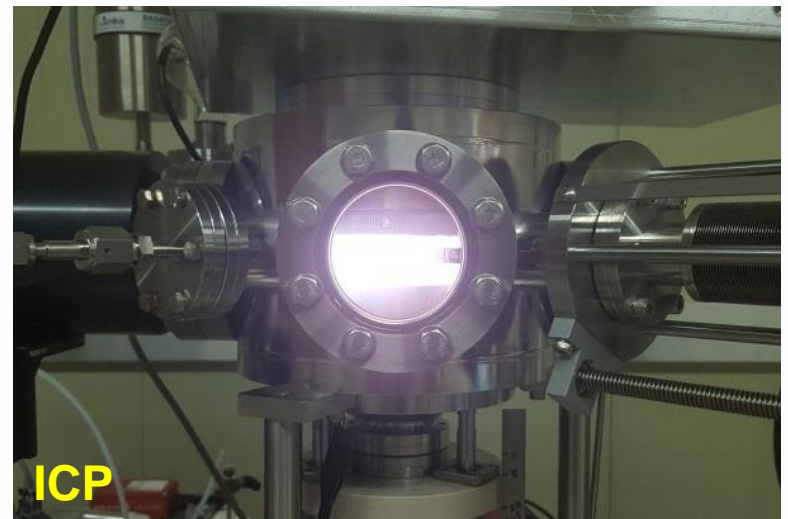
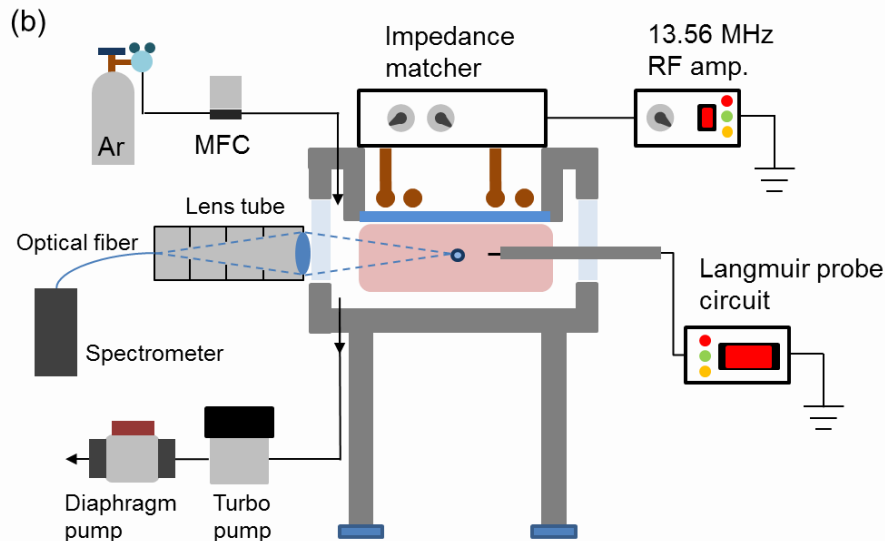
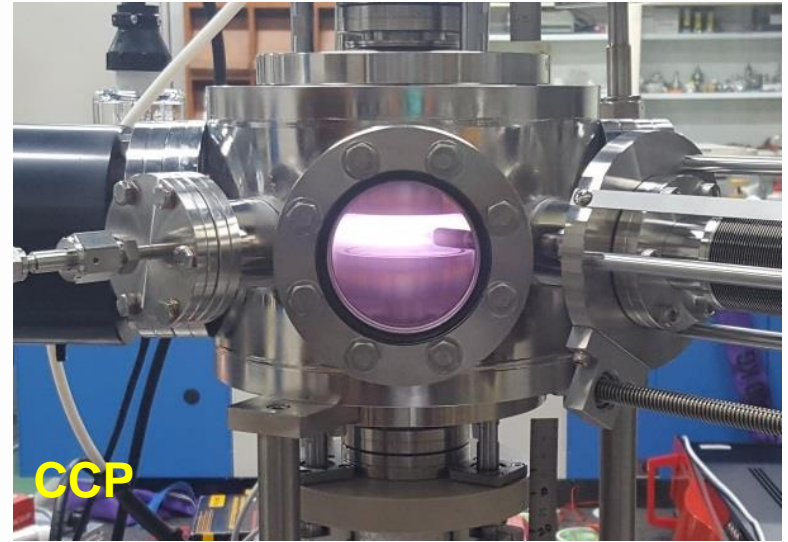
- Collisional-radiative modeling (CRM) for He, H/D plasmas
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OES in low temperature plasmas

Experimental setup



CCP: Capacitively Coupled Plasma
ICP: Inductively Coupled Plasms



CRM for low temperature He I plasma



Steady state balance equation for excited levels

From the particle balance equation $\frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{u}) = \frac{\delta n}{\delta t}$, $n\mathbf{u} = -\nabla(D_a n)$

$$\frac{\partial N_i}{\partial t} - \nabla \cdot (\nabla D_a N_i) = \left(\frac{\delta N_i}{\delta t}\right)_{CR}, \quad \nabla \cdot (\nabla D_a N_i) \approx v_i^d N_i, \quad \frac{\partial N_i}{\partial t} = 0$$

In the weakly ionizing plasma conditions $N_0 \alpha_I \gg n_+ \alpha_R$, $n_+ \approx n_e$

Populating terms

Depopulating terms

$$\rightarrow \sum_{j \neq i} n_e \alpha_{ji}^{ex} N_j + \sum_{j > i} \eta_{ji}(N_i) A_{ji} N_j = \sum_{j \neq i} n_e \alpha_{ij}^{ex} N_i + \sum_{j < i} \eta_{ij}(N_j) A_{ij} N_i + n_e \alpha_i^I N_i + \sum_j \alpha_{ij}^I N_i N_j + v_i^d N_i$$

Nonlinear terms

Ground level population

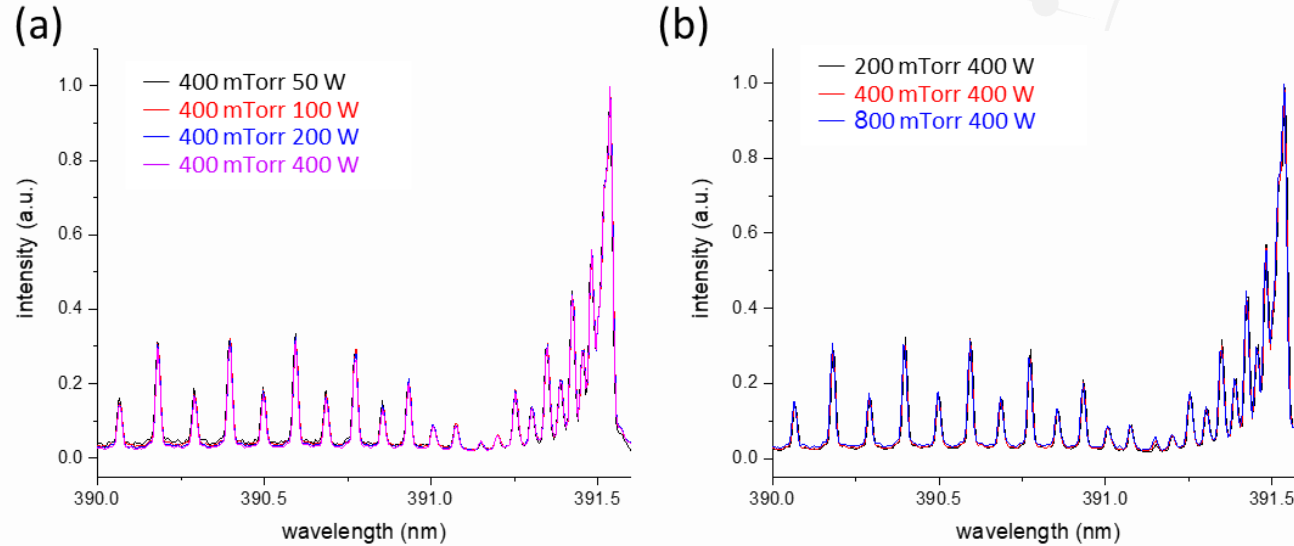
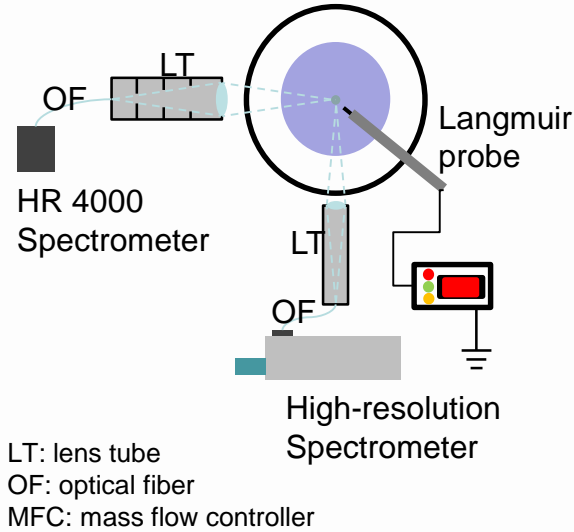
$$p_{tot} = N_0 k_B T_g + n_+ k_B T_i + n_e k_B T_e \approx N_0 k_B T_g$$

$$N_0 \cong \frac{p_{tot}}{k_B T_g} \text{ (constant)}$$

Diagnostics for plasma parameters

Minimization of $\Delta(n_e, T_{eff}, R_{eff}, L_{eff}) = \sum \left(\frac{I_{ik}^{CRM} - I_{ik}^{OES}}{I_{ik}^{OES}} \right)^2$, $I_{ik}^{CRM} = \frac{N_i \eta_{ik} A_{ik}}{\lambda_{ik}}$

Gas temperature for He I plasma



Measured spectra for N_2^+

Gas temperature

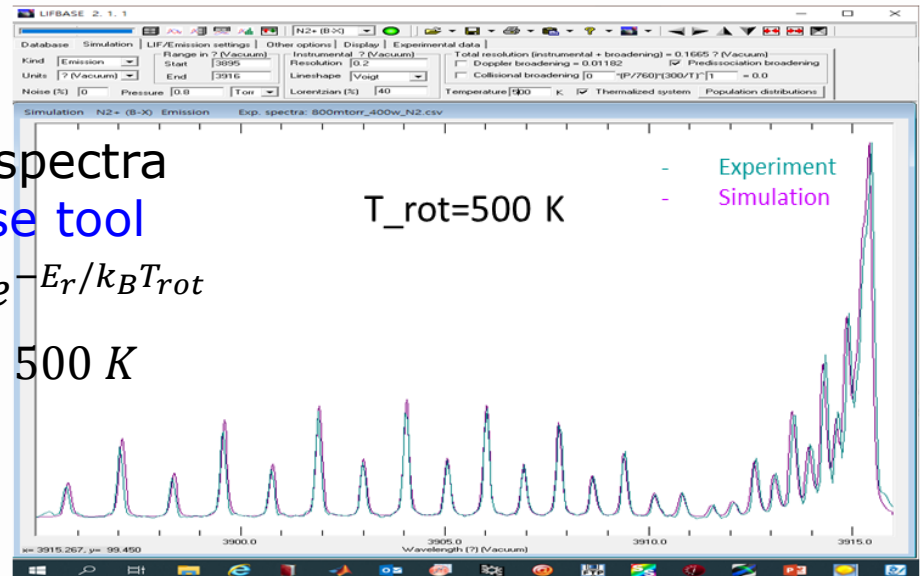
T_g can be estimated with the N_2^+ rotational temperature obtained from the emission spectra of the N_2^+ transition $B^2\Sigma_u^+, v=0 \rightarrow X^2\Sigma_g^+, v'=0$ by inserting small amount of N_2 gas into the He plasma.

(c)

Modeled spectra by LIFbase tool

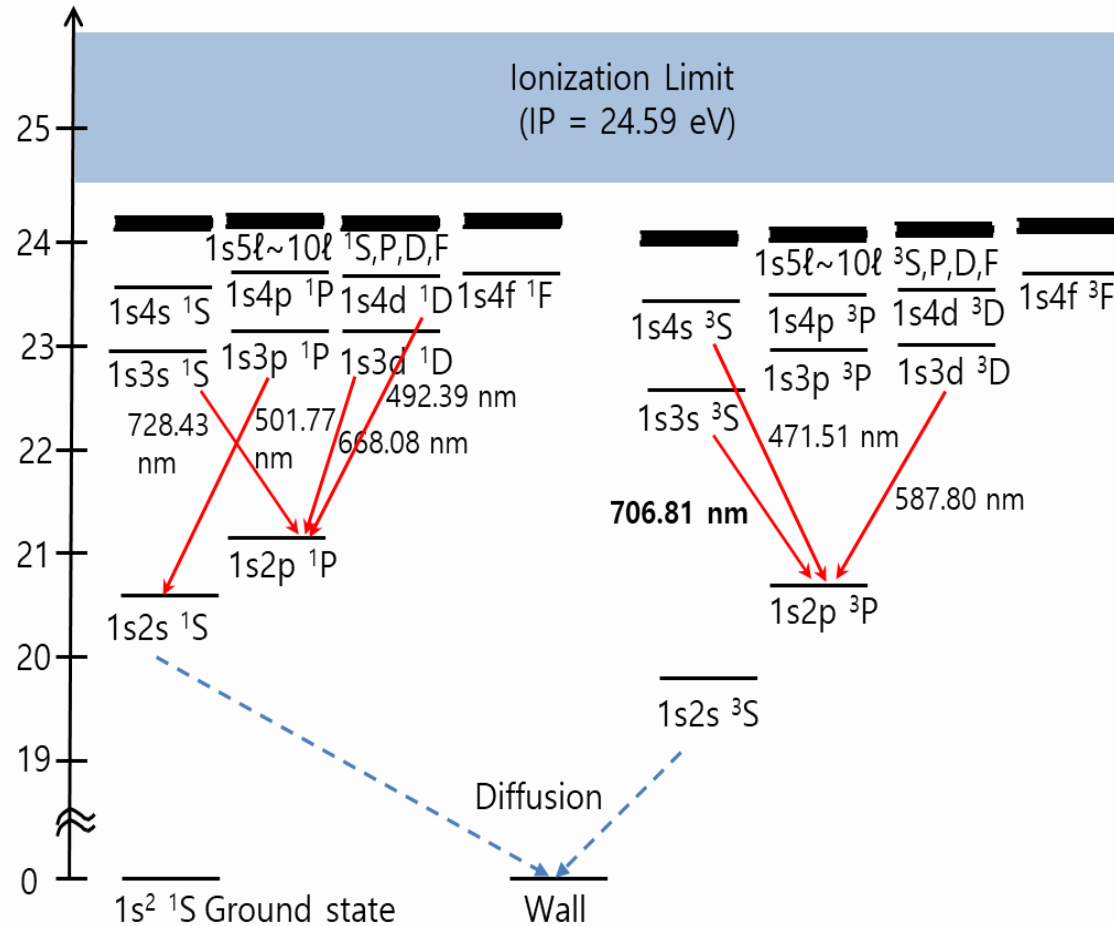
$$I = D_0 k^4 S e^{-E_r/k_B T_{rot}}$$

$$\Rightarrow T_g \approx T_{rot} = 500 \text{ K}$$



Energy levels and Kinetic processes of He I

Energy levels



up to 1snl ($n = 10, \ell = 3$)

Kinetic processes

$$1. \text{He} + e \rightarrow \text{He}^* + e \quad \alpha_{ij}^{ex} \quad [1]$$

$$2. \text{He} + e \rightarrow \text{He}^+ + 2e \quad \alpha_i^I \quad [1]$$

$$3. \text{He}^* \rightarrow \text{He} + h\nu \quad \lambda_{ij}, A_{ij} \quad [2]$$

$$4. \text{He}(1s2\ell) + \text{He}(1s2\ell') \rightarrow \text{He}^+ + \text{He} + e \quad \alpha_{ij}^I$$

$$2.9 \times 10^{-9} (T_g/300)^{1/2} \text{ (cm}^3/\text{s)}$$

$$5. \text{He}(1s2s) \rightarrow \text{to wall} \quad v_i^d$$

$$v_i^d = D_a \left(\left(\frac{2.405}{R_{eff}} \right)^2 + \left(\frac{\pi}{L_{eff}} \right)^2 \right),$$

$$D_a = 8.992 \times 10^{-2} \frac{T_g^{3/2}}{p} \text{ (cm}^3/\text{s)}$$

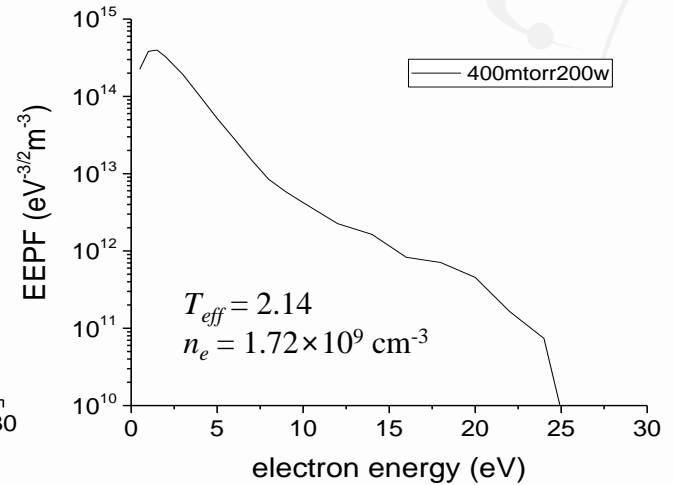
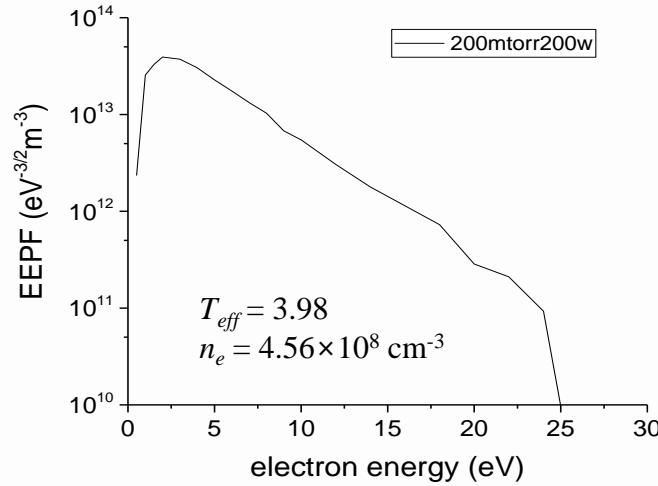
[1] Y. Ralchenko, R. K. Janev, T. Kato, D. V. Fursa, I. Bray, F. J. de Heer, Atomic Data and Nuclear Data Tables 94 (2008) 603.

[2] G.W.F Drake, D.C. Morton, Astrophys. J. Suppl. Series 170 (2007) 251.

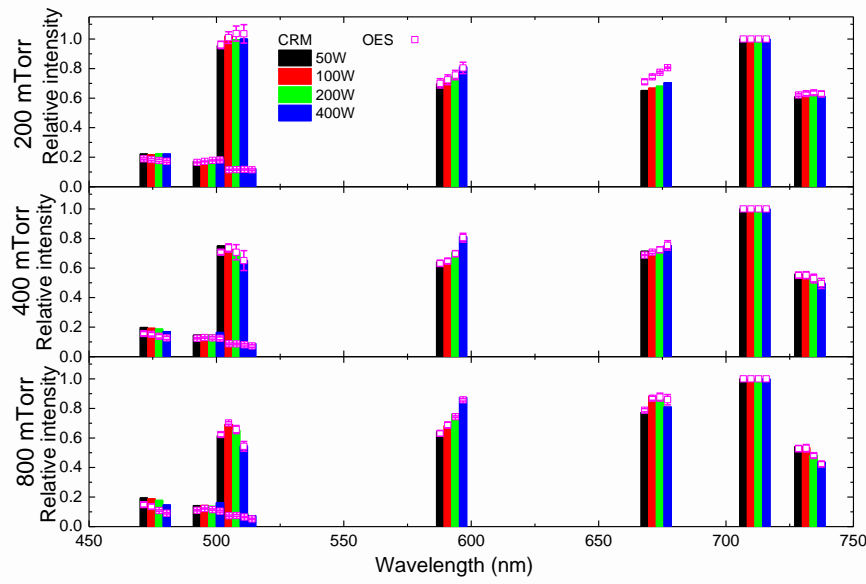
Results for diagnostics



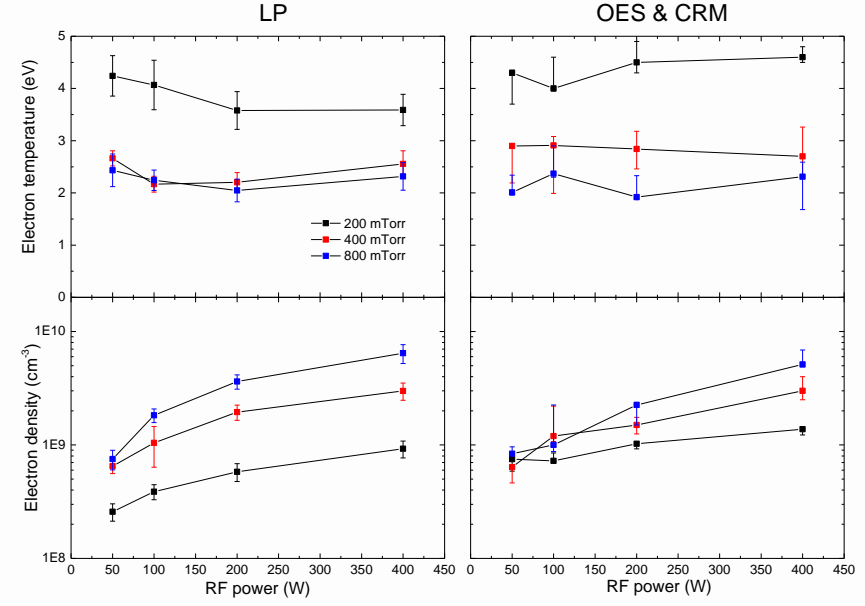
Non-Maxwellian electron energy distribution



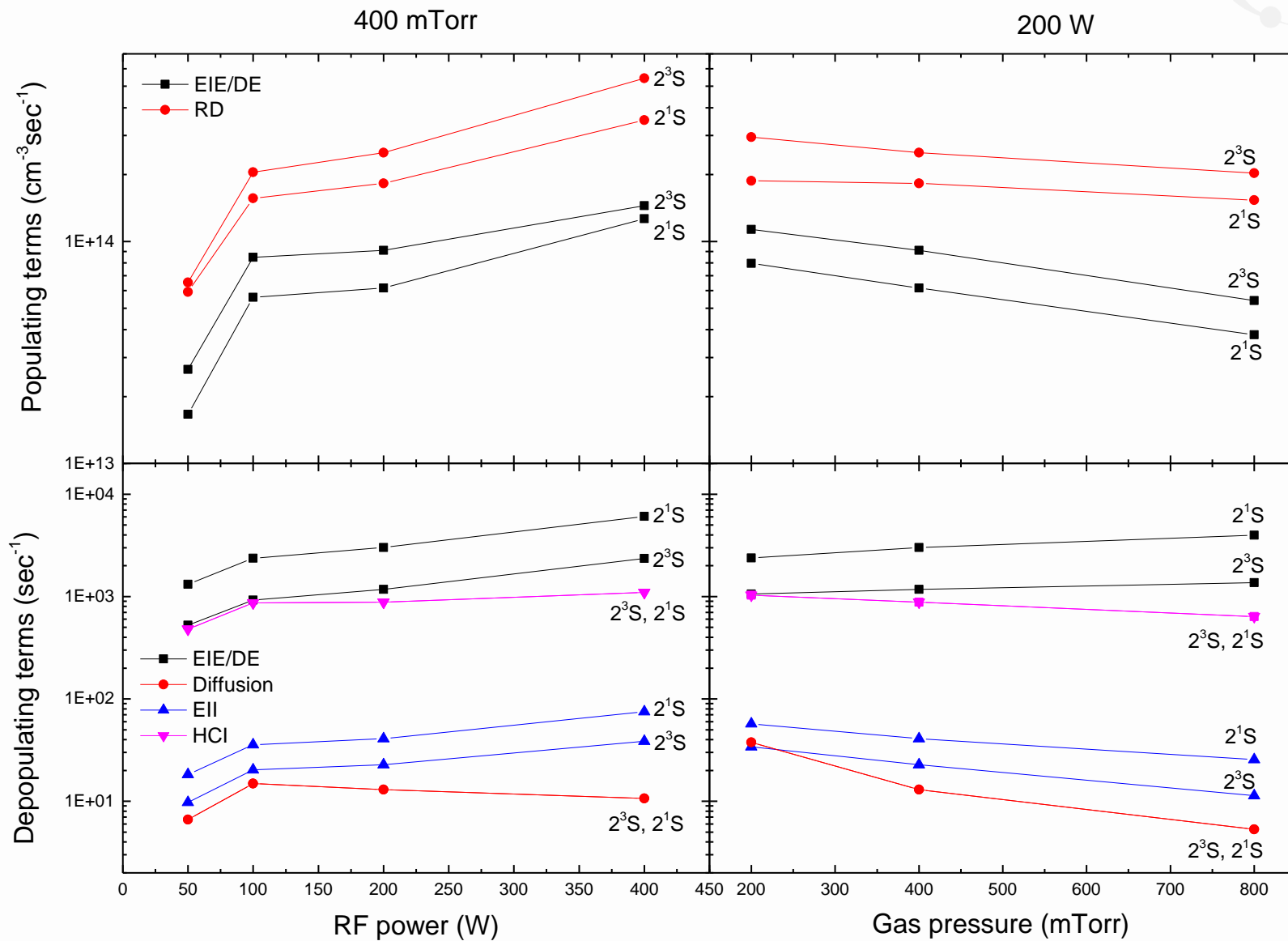
OES and CRM modeling spectra



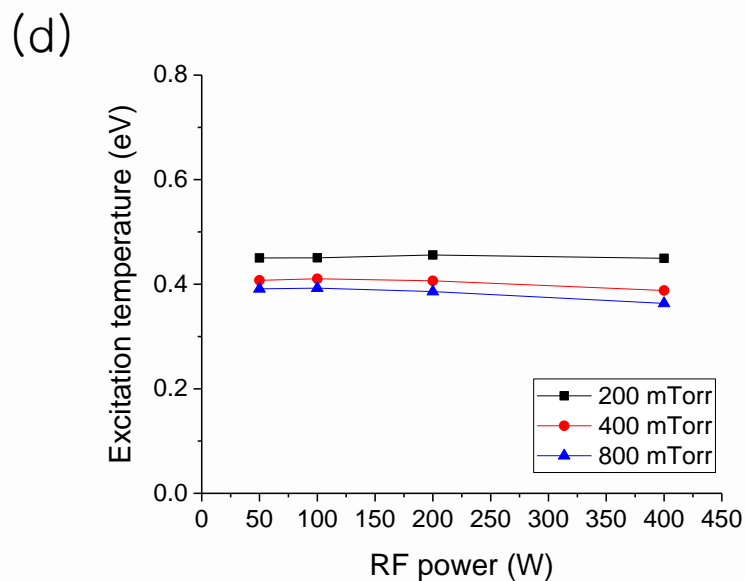
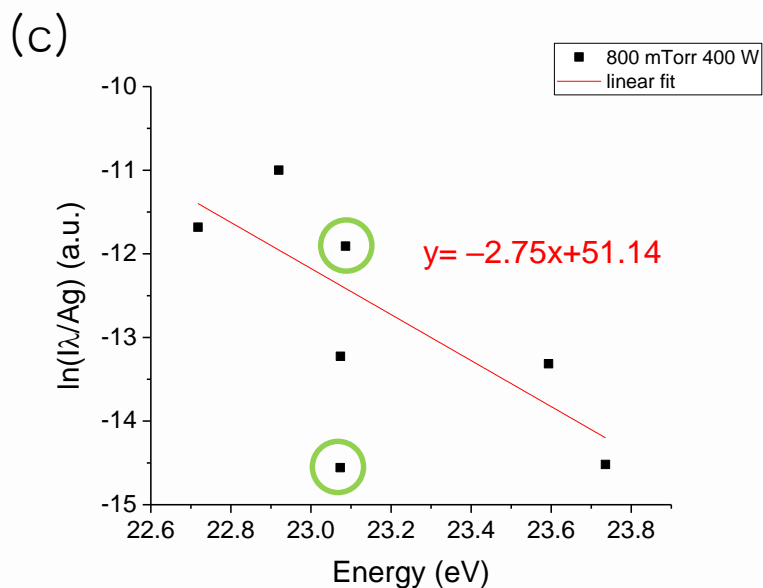
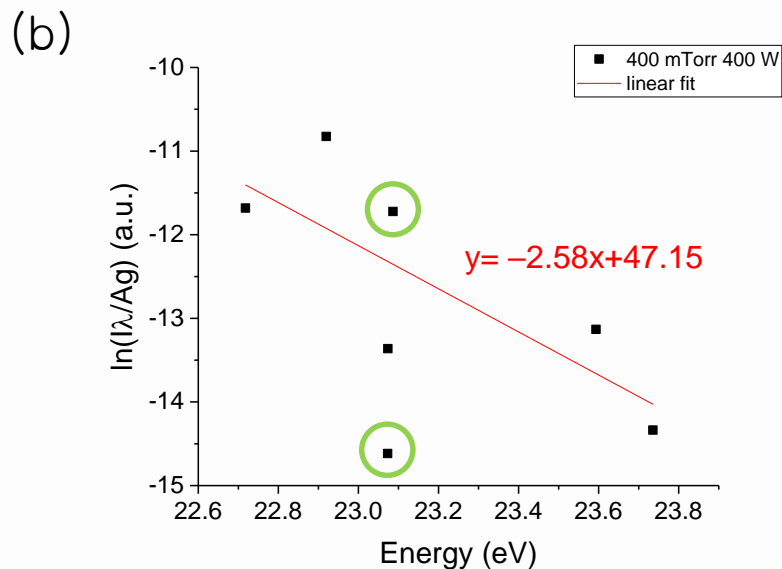
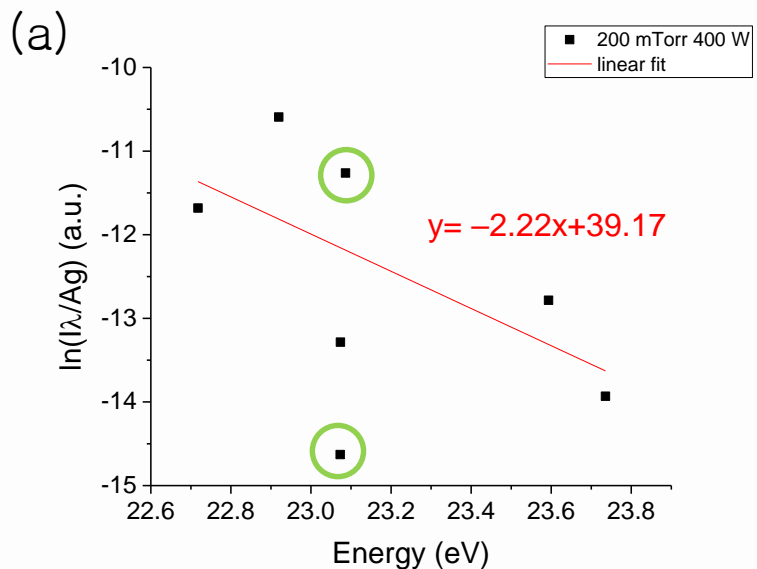
Electron temperature and density



Population kinetics for He I



Deviation from Boltzman distributions



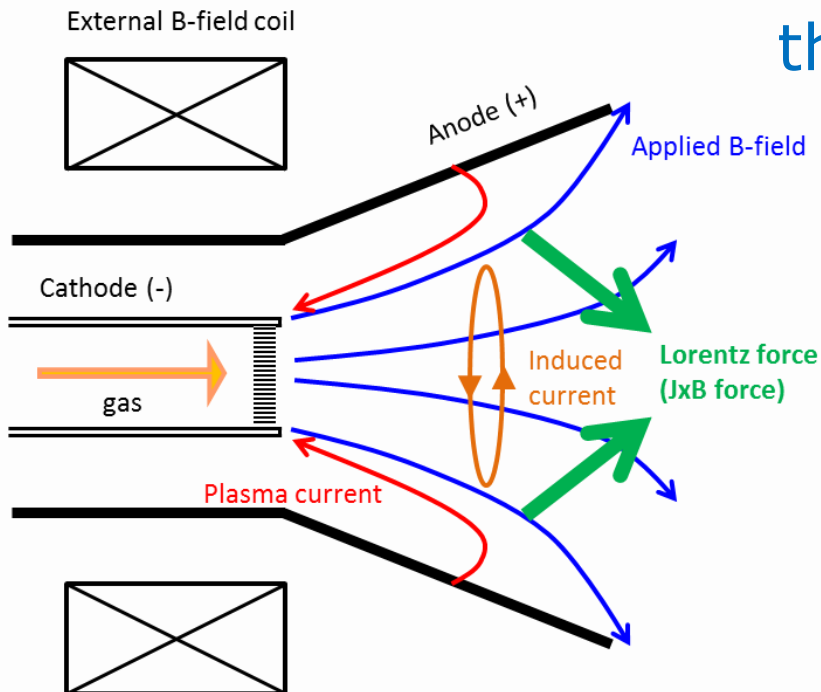
KAERI Plasma Beam Irradiation Facility



Motivation of the construction

In order to develop divertor materials and cooling techniques resisting **high heat** and **particle fluxes** (heat flux of 10 MW/m^2 and particle flux of $10^{24} /\text{m}^2\text{s}$ will come in ITER and much larger heat and particle fluxes will come in DEMO), we have constructed **lab-scale** divertor plasma simulator

Applied field-magnetoplasdynamic (AF-MPD) thruster concept

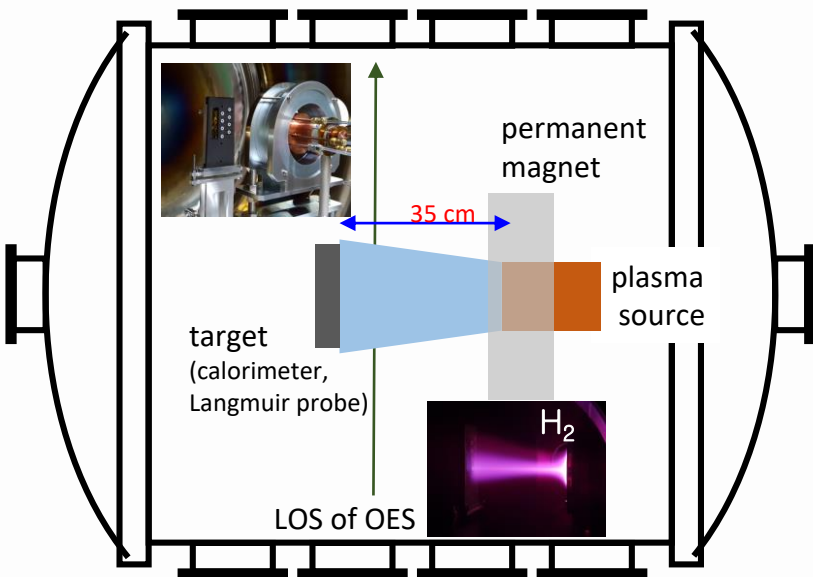


type I/ type II

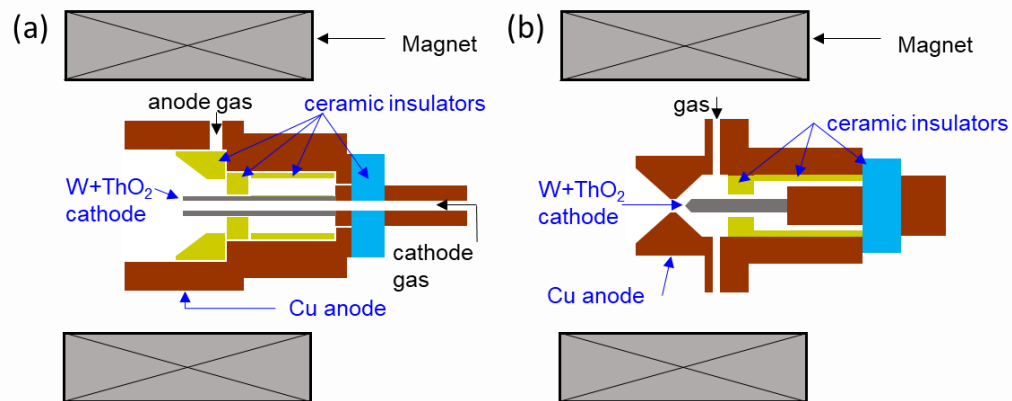
type I/ type II

- Anode radius = $4/2 \text{ cm}$, cathode radius = $0.6/0.4 \text{ cm}$
- Anode material: Cu, cathode material: W+ThO₂ (2%)
- Insulating material : ceramic (Al₂O₃)
- Sustain power supply : DC 10-20 kW
- External B-field: 0.17 T (NdFeB permanent magnet)
- Both anode & cathode can be water-cooled

KPBIF schemes and measurement



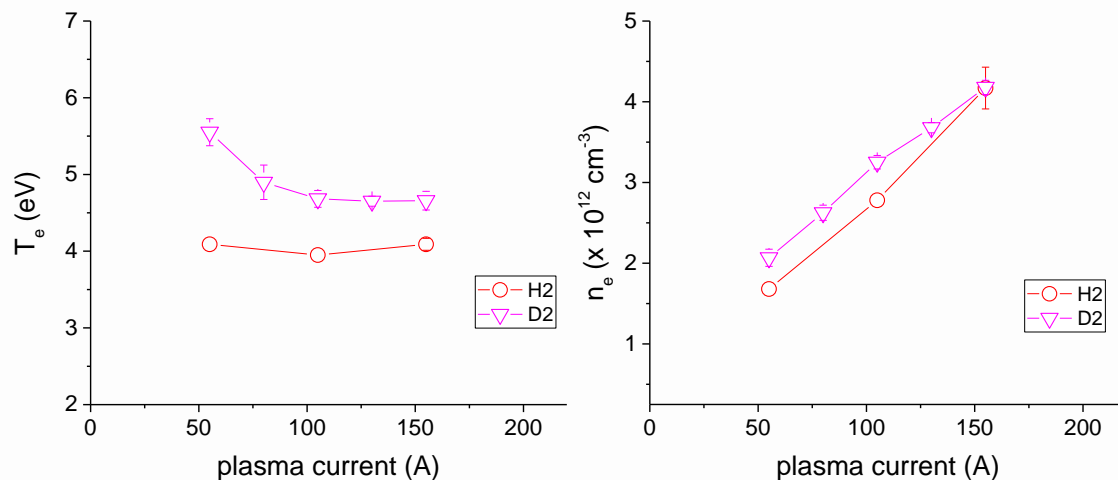
Plasma source schemes



Type I for Ar

Type II for H₂, D₂, He

Electron temperature and density by LP



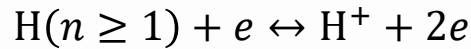
CRM for H/D plasma



Considered processes



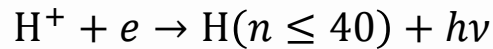
(1) [1] [1] R. K. Janev, D. Reiter and U Samm, Collision processes in low-temperature hydrogen plasmas, Report JUEL-4105 (2003).



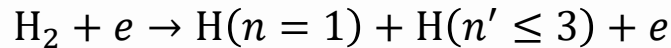
(2) [1] [2] R. K Janev et al., Elementary processes in Hydrogen-Helium plasmas, (Berlin: Springer 1987)



(3) [1] [3] W. L. Wiese, and J. R. Fuhr, Accurate atomic transition probabilities for H, He, Li, J. Phys.. Chem. Ref. Data **38** 565 (2009).



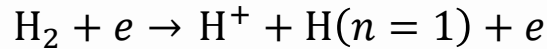
(4) [1] [4] P. del Mazo-Sevillano, ..., D.-H. Kwon, O. Roncero, Molecular physics e2183071



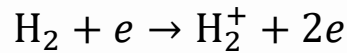
(5) [2] ❖ The cross sections for electron collisions and radiative transitions of D species were used by those of H species. The heavy particle collision cross section of D species was from the ab-initio calculation for D [4]. The mass effect for rate coefficients and mobility of D were taken into account.



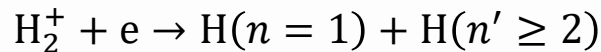
(6) [2]



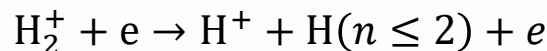
(7) [1]



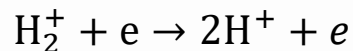
(8) [1]



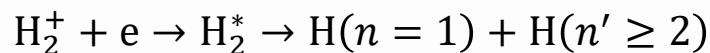
(9) [3]



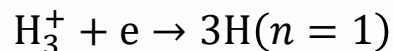
(10) [3]



(11) [3]



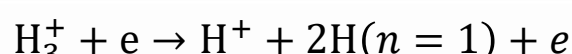
(12) [1]



(13) [3]



(14) [3]



(15) [3]



(16) [4]

$$\alpha(T_{12}) = \frac{4}{\sqrt{\pi}v_{T_{12}}^3} \int_0^\infty \sigma(v_{12}) \exp\left(-\left(v_{12}/v_{T_{12}}\right)^2\right) v_{12}^3 dv_{12}$$

$$T_{12} = (m_2 T_1 + m_1 T_2) / (m_1 + m_2)$$

$$v_{T_{12}} = \sqrt{2(m_1 + m_2)T_{12} / m_1 m_2}$$

CRM for H/D plasma



For atomic levels n_i ($i = 1 - 40$)

$$D_{AH^+} = T_e K_1^0 \left(\frac{760 T_m}{p \cdot 273} \right)$$

$$p = n_{H_2} T_m$$

$$K_1^0 = 15.9 (H^+), 11.2 (D^+) \\ (cm^2 V^{-1} s^{-1})$$

$$D_{AH^+} : D_{AH_2^+} : D_{AH_3^+} = 1 : \frac{\sqrt{2}}{\sqrt{3}} : \frac{\sqrt{5}}{3}$$

$$\frac{dn_i}{dt} = \sum_{j>i}^{40} \eta_{ji} A_{ji} n_j - \left(\sum_{j<i} \eta_{ij} A_{ij} + \frac{Q}{V} + \frac{\gamma}{\tau} \delta_{i1} \right) n_i + \\ n_e \left(\sum_{j \neq i} \beta_{1,ji} n_j - \sum_{j \neq i} \beta_{1,ij} n_i - \beta_{2i} n_i + \beta_{4i} n_{H^+} \right) + \\ n_e (\beta_{5i} + \beta_6 \delta_{i2} + \beta_7 \delta_{i1}) n_{41} + n_e (\beta_{9i} + \beta_{10i} + \beta_{12i}) n_{42} + \\ n_e (\beta_{13} \delta_{i1} + \beta_{14} \delta_{i2} + \beta_{15} \delta_{i1}) n_{43} + n_{41} \beta_{16} \delta_{i1} n_{42} + \\ \left(\sum_{j=42}^{43} \varsigma_{aj} \left(\frac{\mu}{R} \right)^2 D_{Aj} n_j + \varsigma_{aD^+} \left(\frac{\mu}{R} \right)^2 D_{AH^+} n_{H^+} \right) \delta_{i1}$$

$$\tau = 2R/v_{th} \\ v_{th} = 2 \sqrt{\frac{2T_a}{\pi M_H}}$$

$$\mu = 2.405$$

For molecule and ions n_i ($i=41,42,43$ for H_2, H_2^+, H_3^+)

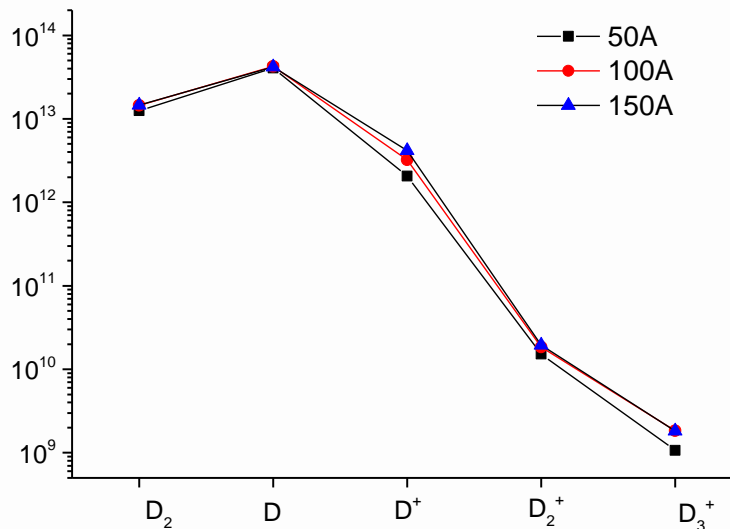
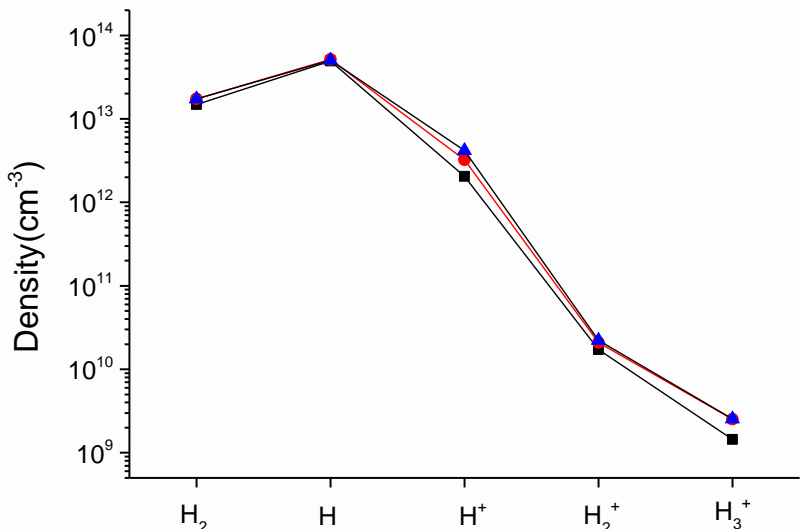
$$\frac{dn_i}{dt} = \delta_{i41} \left(n_e \beta_{14} n_{43} + \frac{Q_{in}}{V} \times 4.48 \times 10^{17} + \frac{\gamma'}{\tau} n_1 + \sum_{j=42}^{43} \varsigma_{mj} \left(\frac{\mu}{R} \right)^2 D_{Aj} n_j + \varsigma_{mH^+} \left(\frac{\mu}{R} \right)^2 D_{AH^+} n_{H^+} \right) \\ - \delta_{i41} n_e (\beta_5 + \beta_6 + \beta_7 + \beta_8) n_i - \delta_{i42} n_e (\beta_9 + \beta_{10} + \beta_{11} + \beta_{12}) n_i - \\ \delta_{i43} n_e (\beta_{13} + \beta_{14} + \beta_{15}) n_i + \delta_{i42} (n_e \beta_8 - n_{41} \beta_{16}) n_i + \\ n_{41} \beta_{16} \delta_{i43} n_{42} - \delta_{i41} \beta_{16} n_{42} n_i - \frac{Q}{V} n_i - (1 - \delta_{i41}) \left(\frac{\mu}{R} \right)^2 D_{Ai} n_i$$

Quasi neutrality condition for H^+ ion n_{H^+} : $n_e = n_{H^+} + n_{H_2^+} + n_{H_3^+}$

rather than **pressure balance equation** :

$$p_{tot} = n_m k_B T_m + k_B T_a \left(\sum_{j=1}^{40} n_j + n_{H^+} \right) + n_e k_B T_e + n_m k_B \left(n_{H_2^+} + n_{H_3^+} \right)$$

Results for H/D CRM



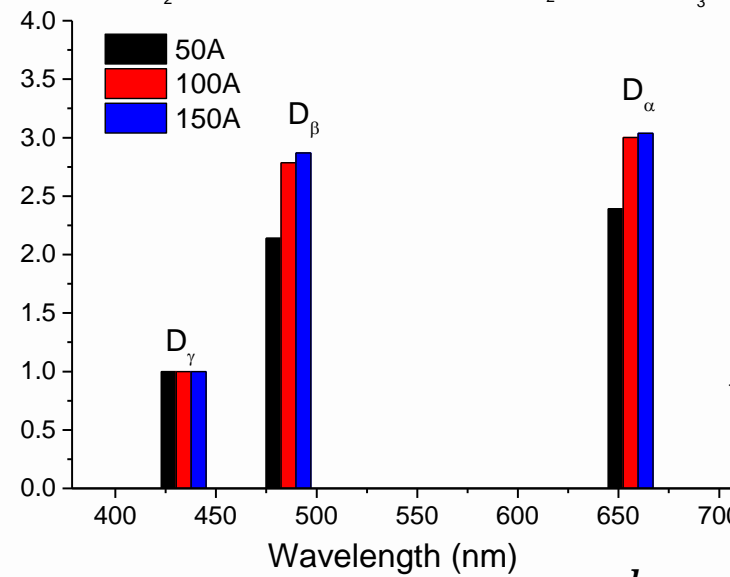
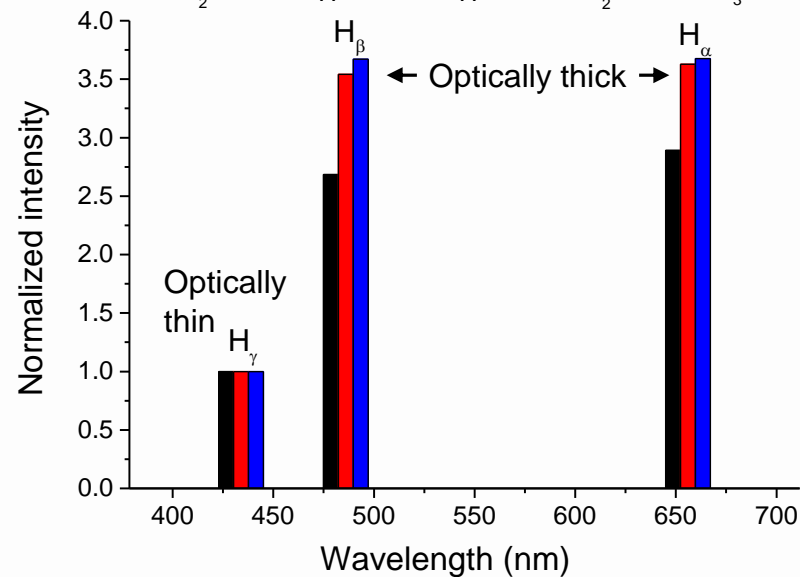
$T_m = 300\text{K}$,
 $T_a = 600\text{K}$,
 $R_{eff} = 40\text{ cm}$,

$T_e = 5.55\text{ (50A)}$,
 4.68 (100A) ,
 4.66 (150A) eV

$n_e = 2.07\text{ (50A)}$,
 3.25 (100A) ,
 4.18 (150A) ,
 $\times 10^{12}\text{ (cm}^{-3}\text{)}$

$Q_{in} = 600\text{ sccm}$,
 $Q = 4800\text{ lps}$,

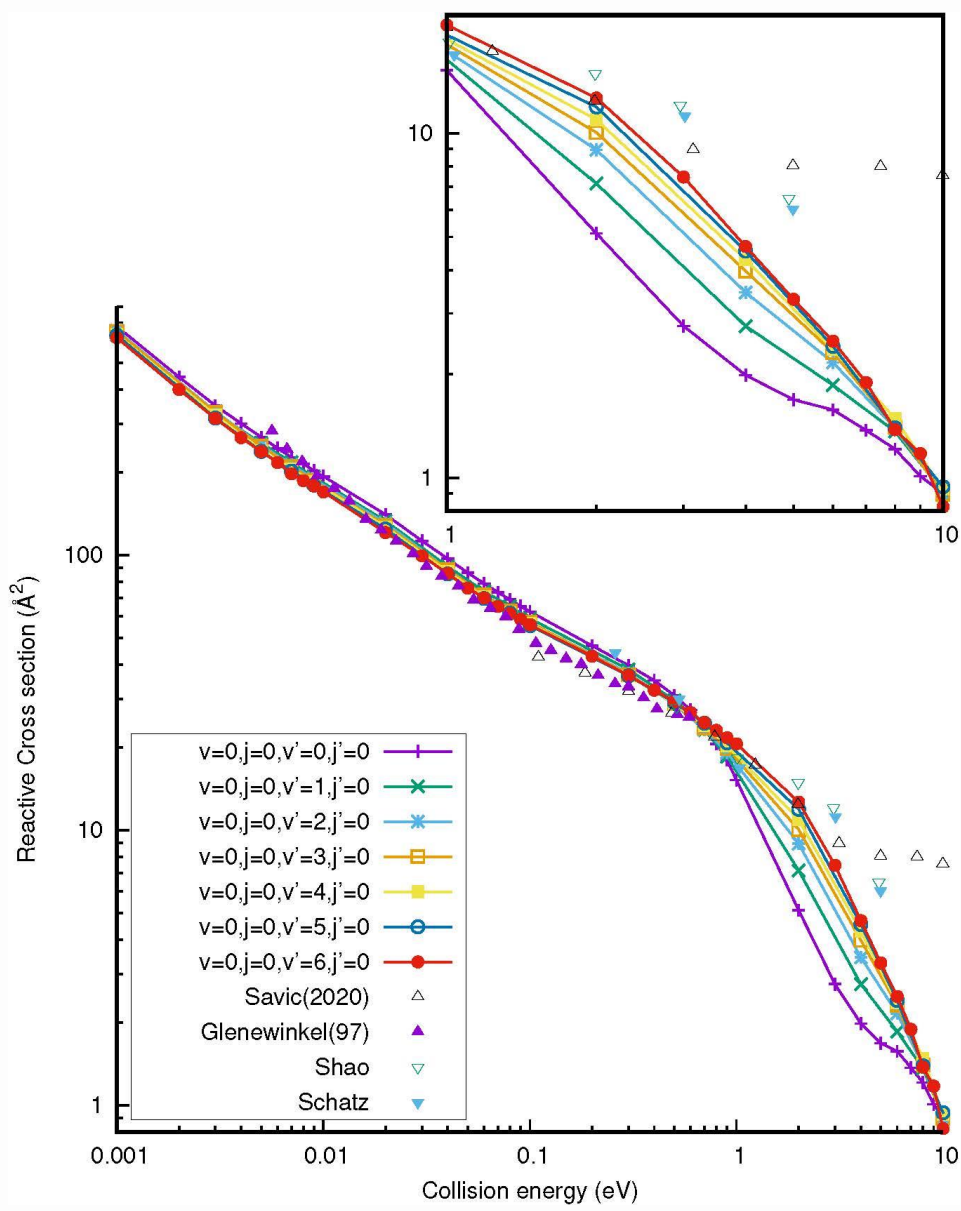
$V = 2.64 \times 10^6\text{ cm}^3$,
 $R = 75\text{ cm}$



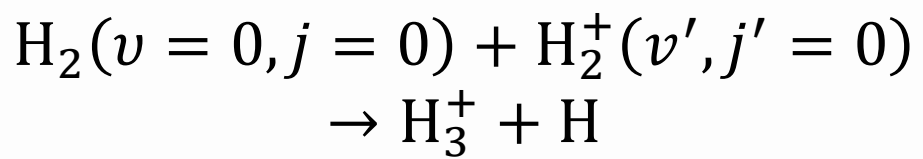
$\text{H} \rightarrow \text{D}, \quad k_0 \uparrow \eta \downarrow$

$$k_0 = \frac{\lambda^3 N_i g_j}{8\pi g_i} \sqrt{\frac{A_{ji} M}{2\pi k_B T_g}}$$

Sensitivity to used atomic data ($X_2 + X_2^+$, $X = H, D$)



Recent data for

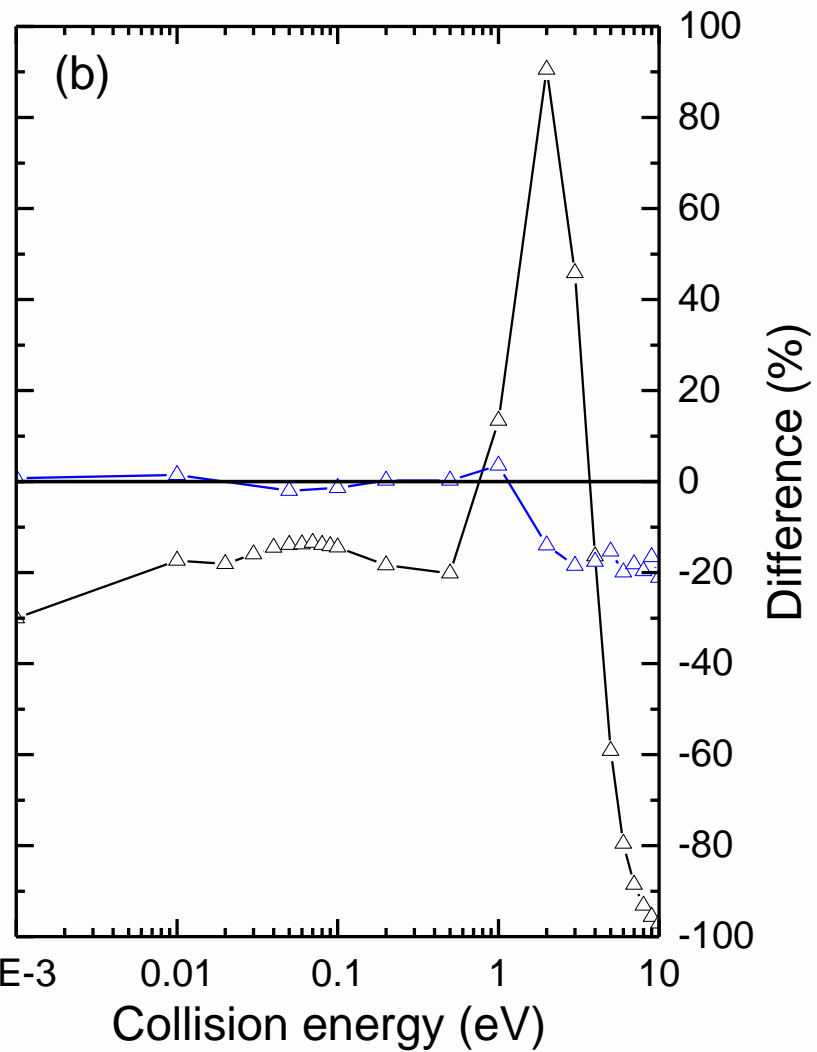
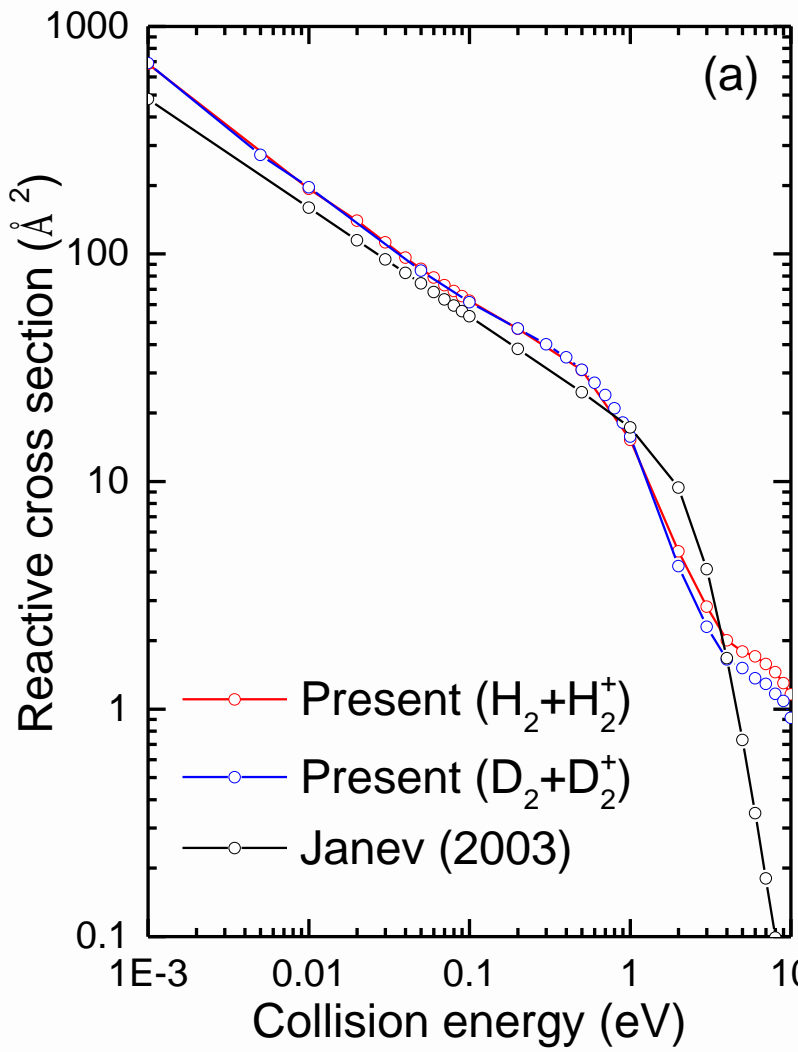


reactive cross sections obtained with QCT (quasi-classical trajectory) calculations

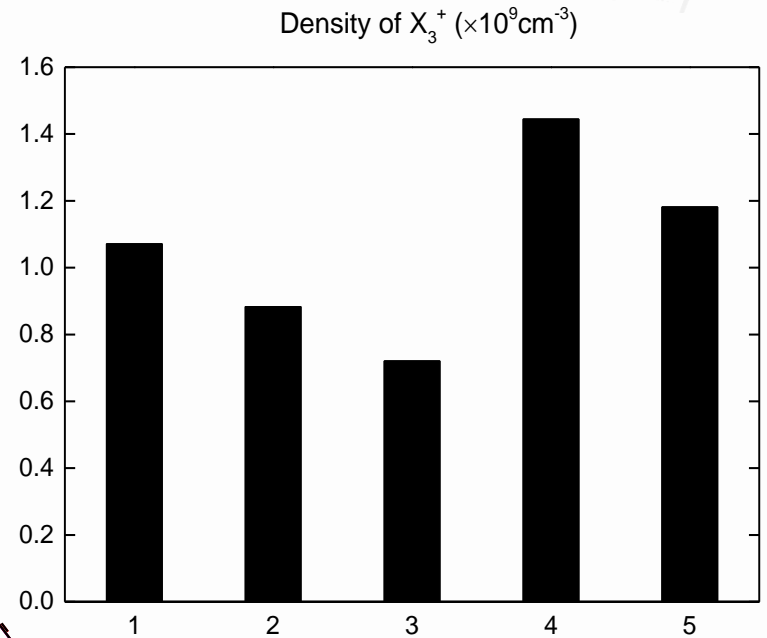
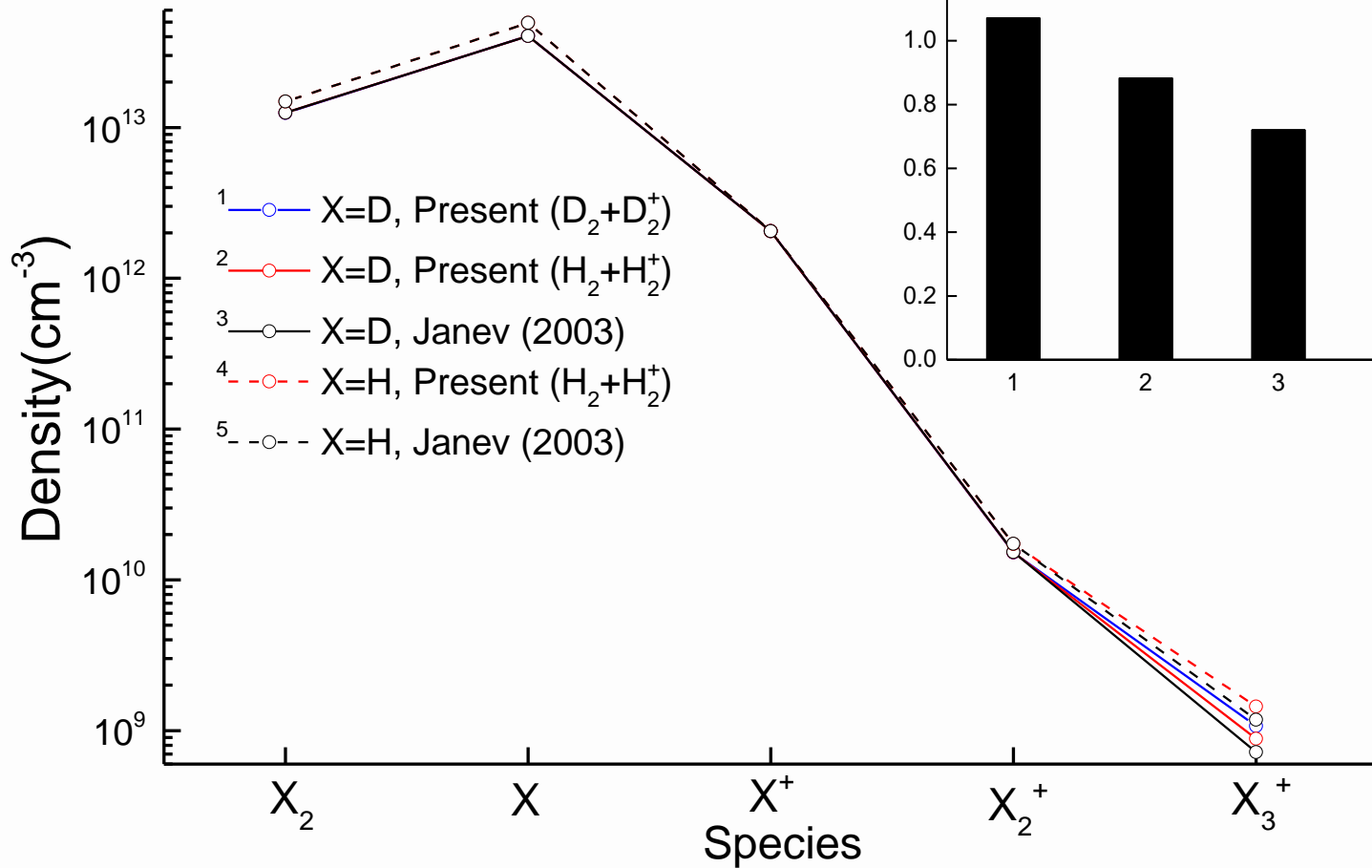
P. del Mazo-Sevillano, ..., D.-H. Kwon, O. Roncero, "Vibrational, non-adiabatic and isotopic effects in the dynamics of the $\text{H}_2 + \text{H}_2^+ \rightarrow \text{H}_3^+ + \text{H}$ reaction: application to plasma modeling", *Molecular physics* e2183071 (2023)

(<https://doi.org/10.1080/00268976.2023.2183071>)

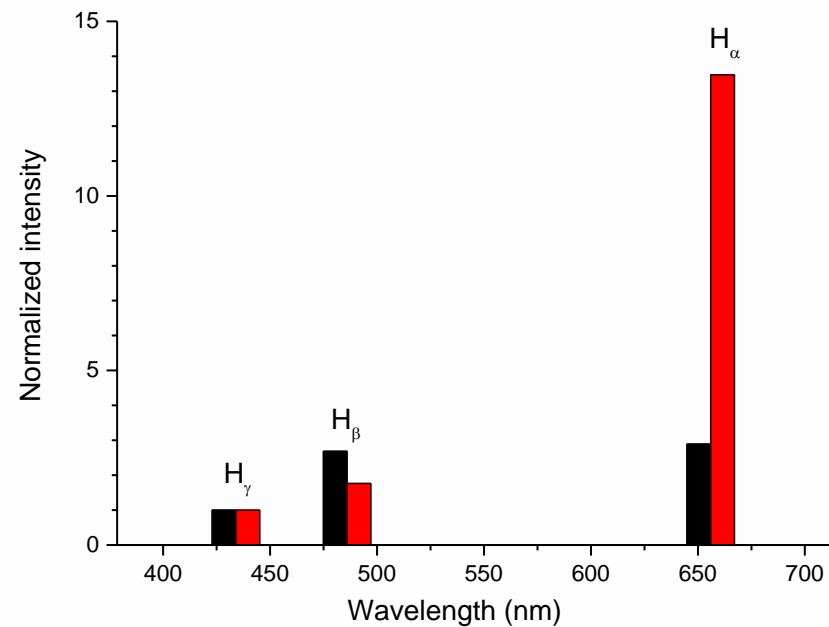
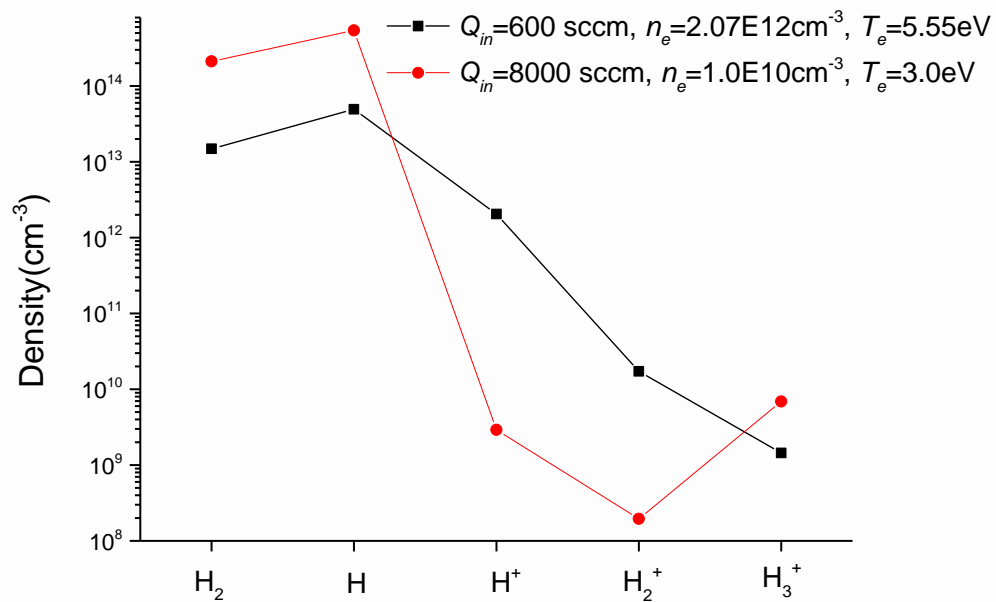
Comparisons of used atomic data



Resulting populations

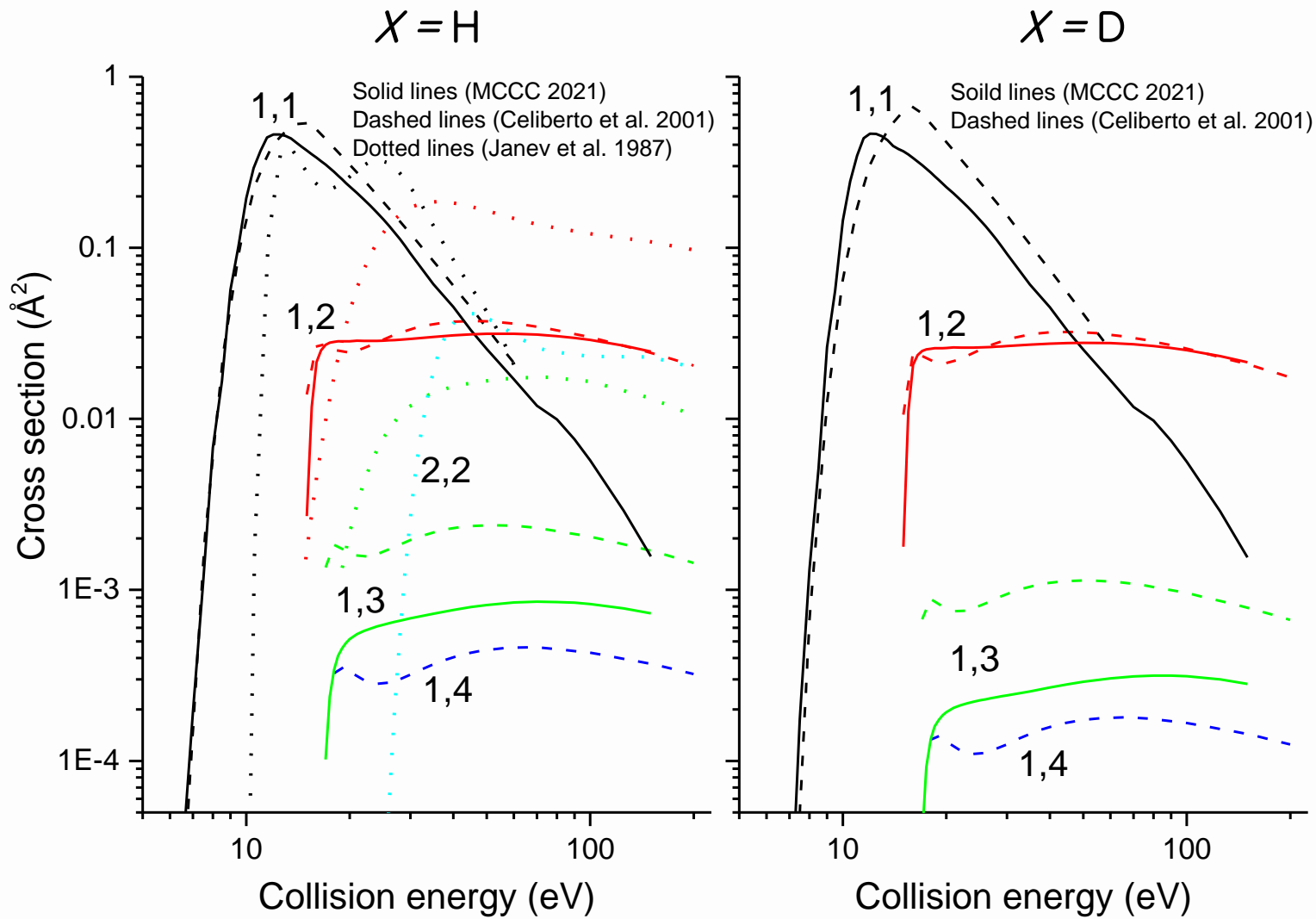
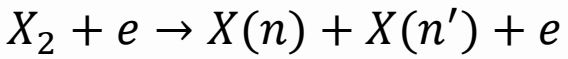


H₃⁺ dominant case





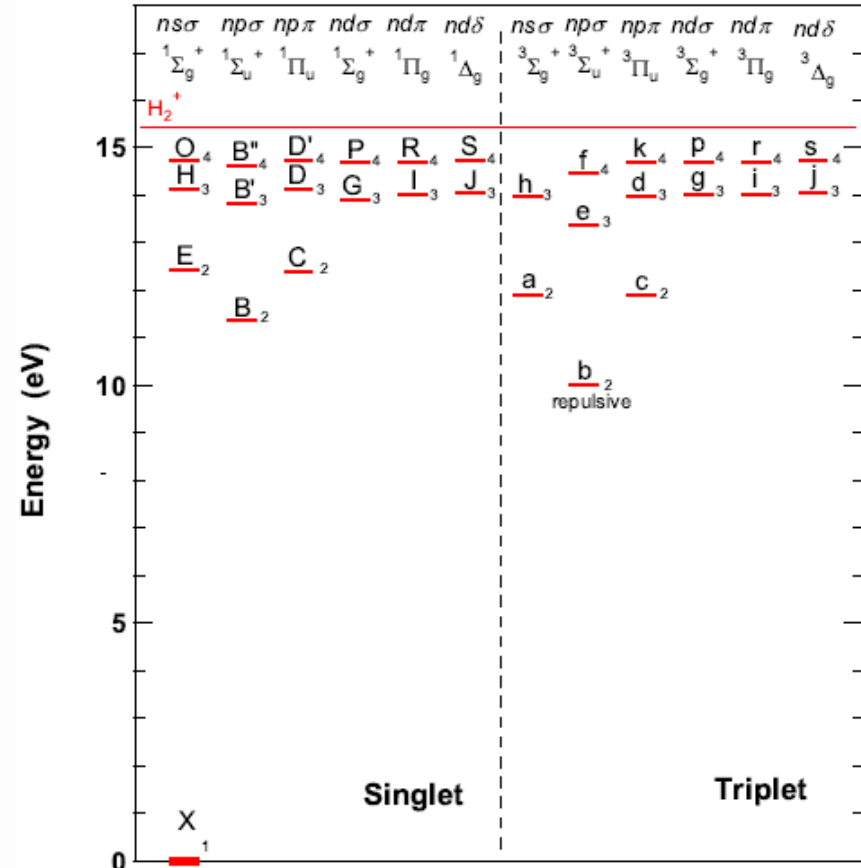
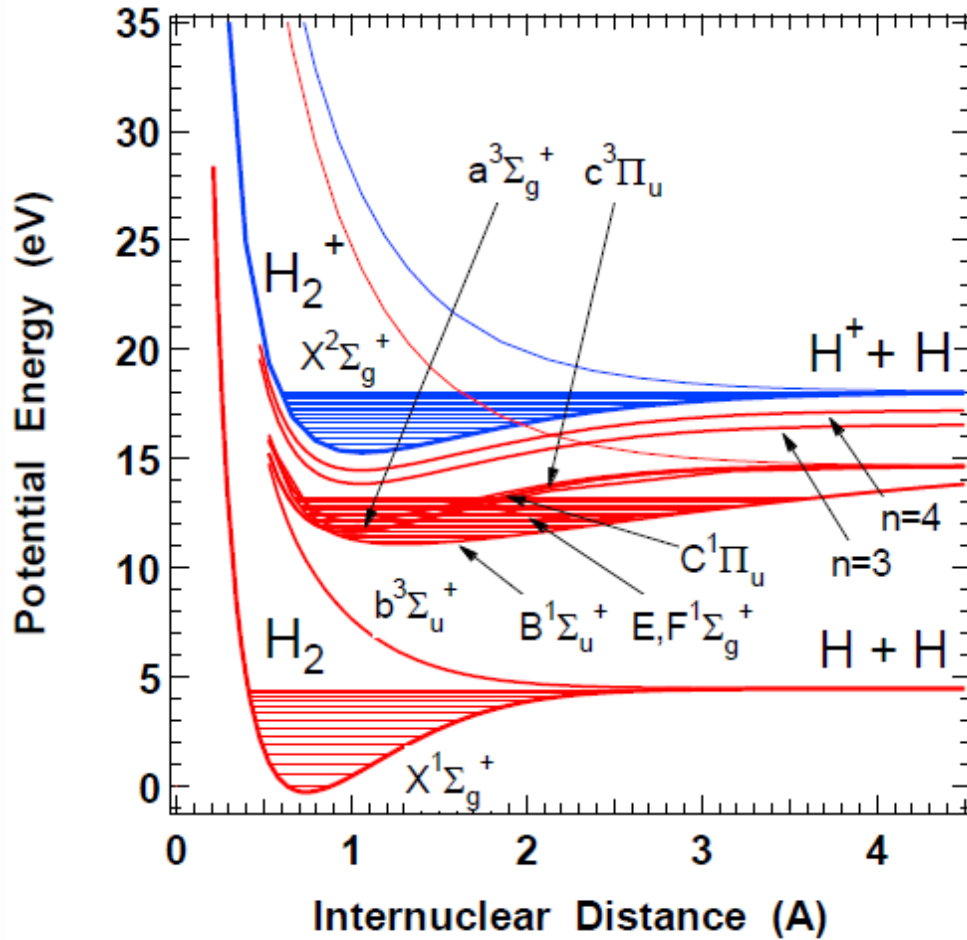
Sensitivity to used atomic data ($X_2 + e, X = H, D$)



R. Celiberto et al.,
Atom. Data Nucl.
Data Tables **77** 161
(2001).

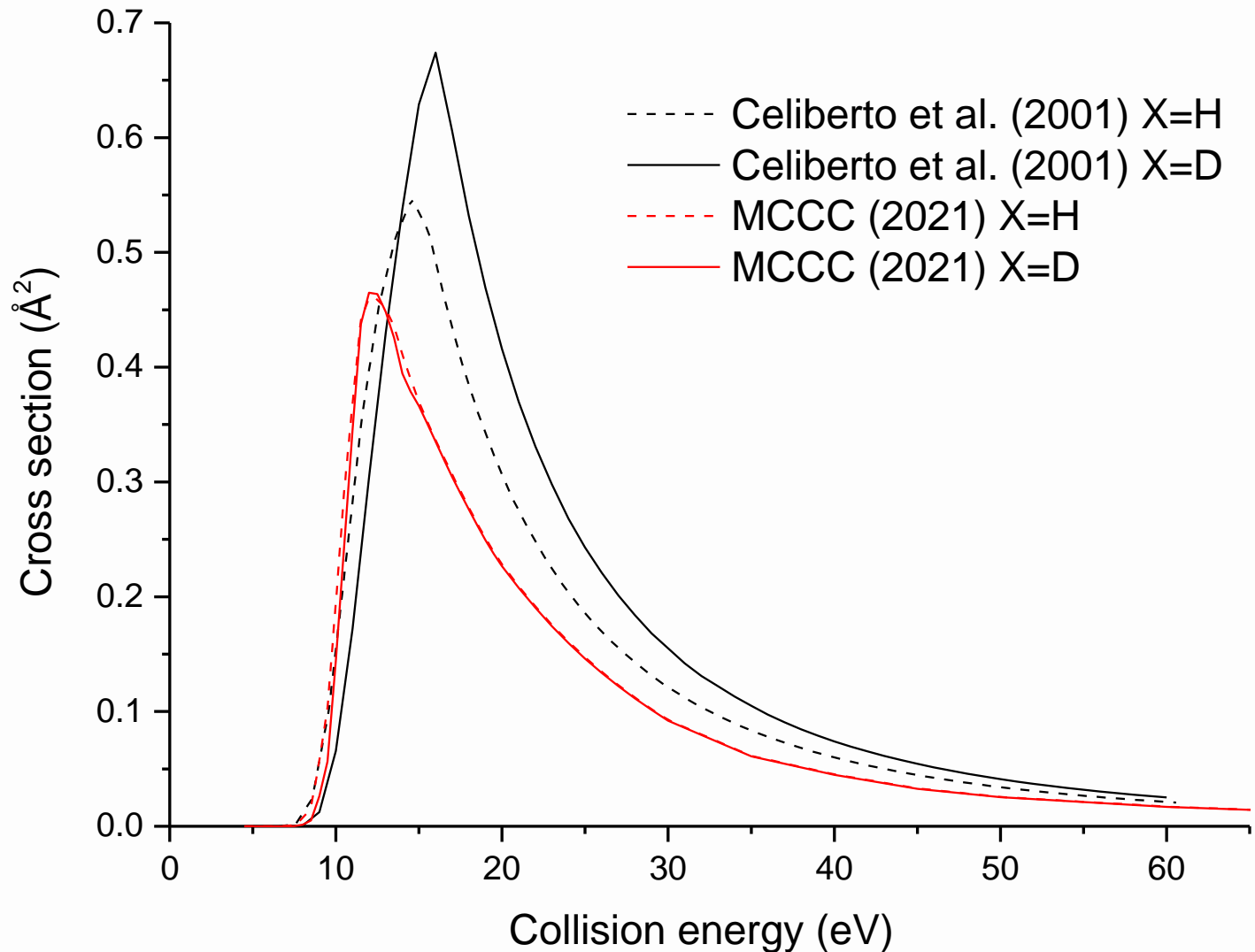
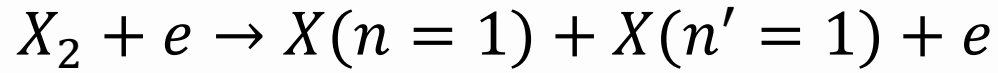
L. H. Scarlett, D. Fursa
et al., Atom. Data
Nucl. Data Tables, **137**
101361 (2021).

H molecular energy levels

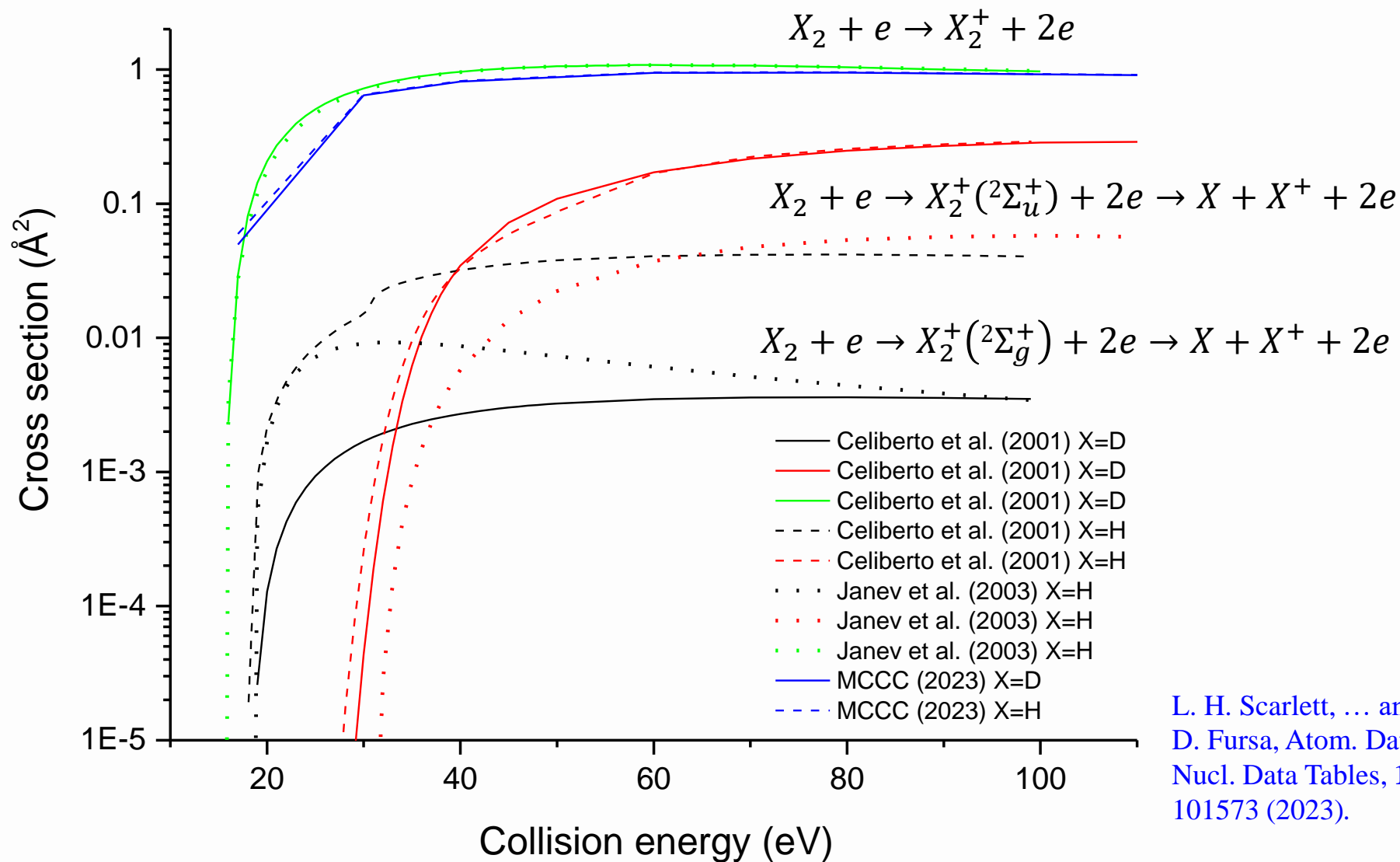


K. Sawada and M. Goto, *Atoms* 4 29 (2016)

Comparisons of used atomic data

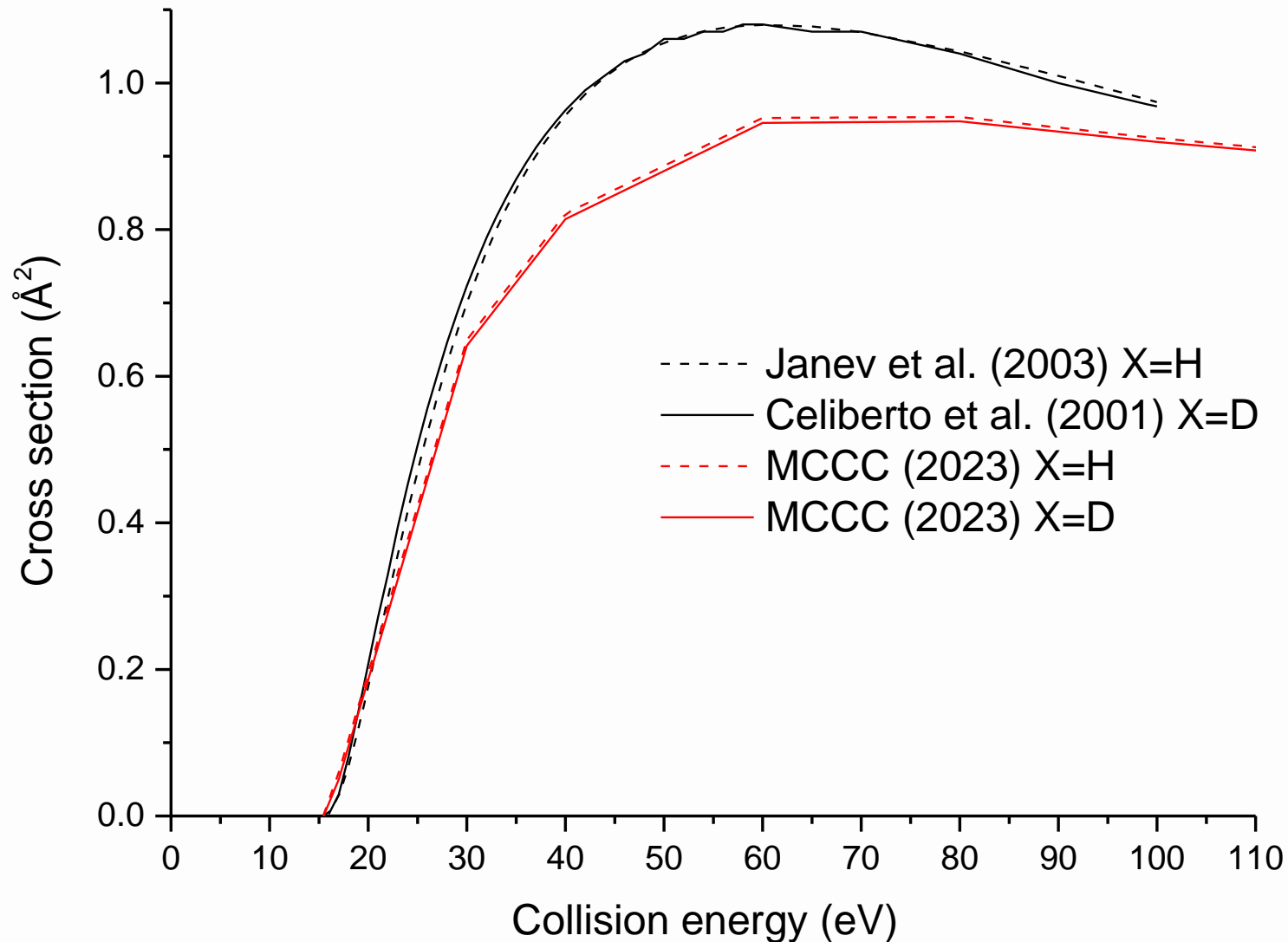
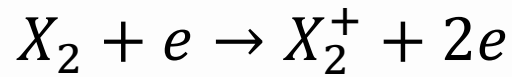


Comparisons of used atomic data

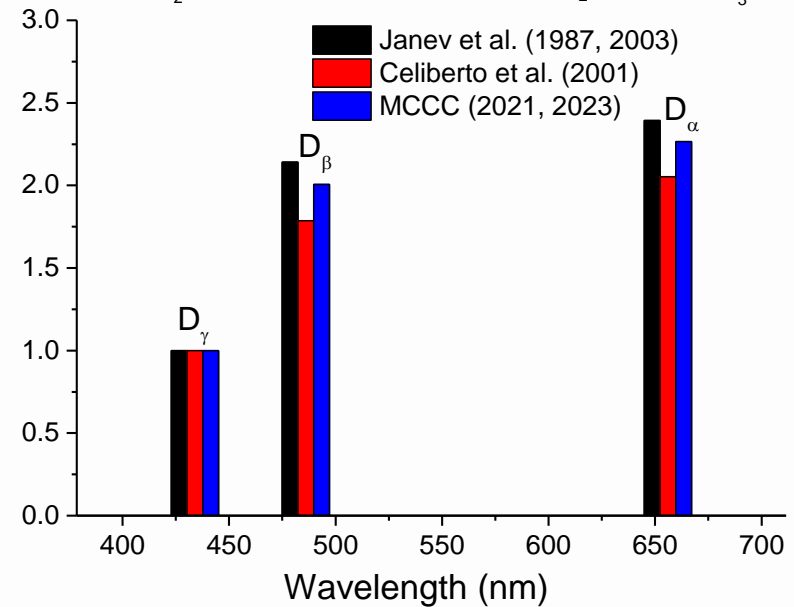
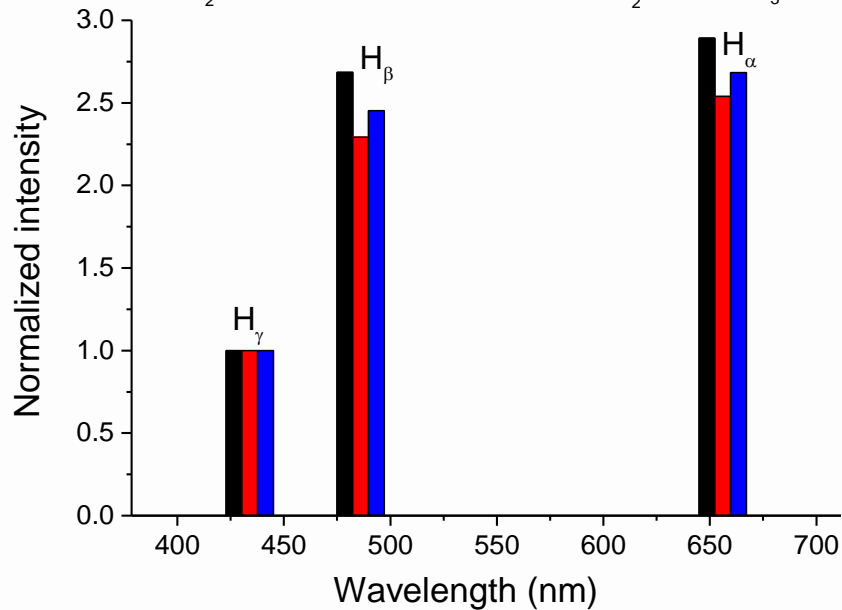
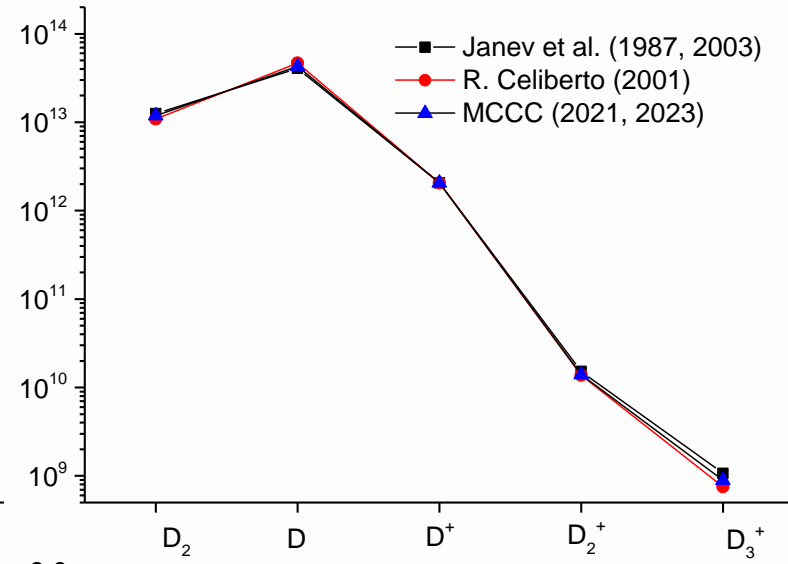
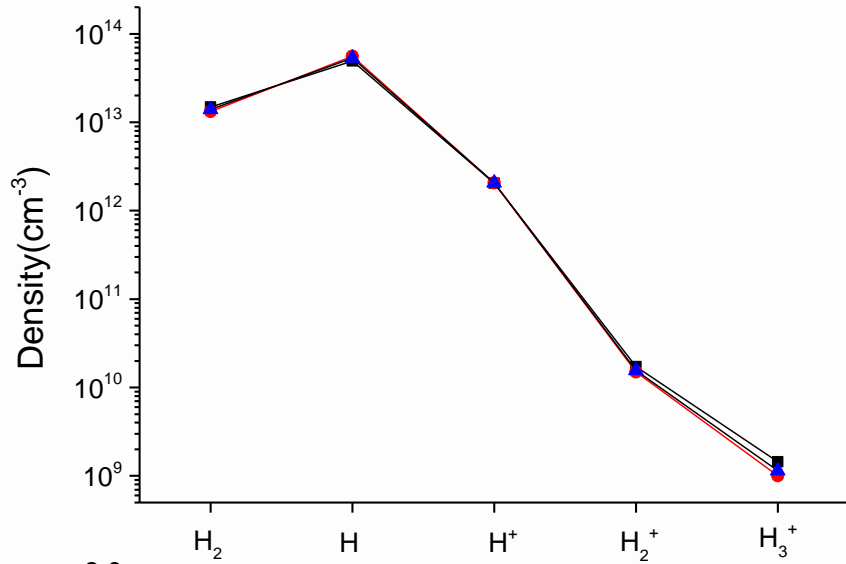


L. H. Scarlett, ... and
 D. Fursa, *Atom. Data
 Nucl. Data Tables*, **151**
 101573 (2023).

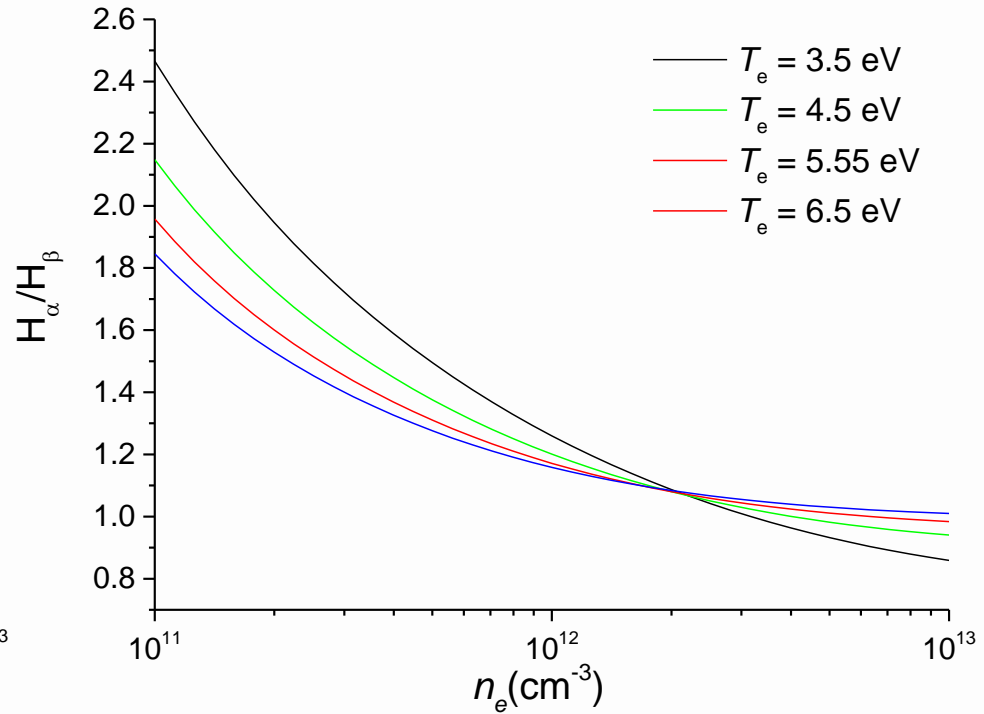
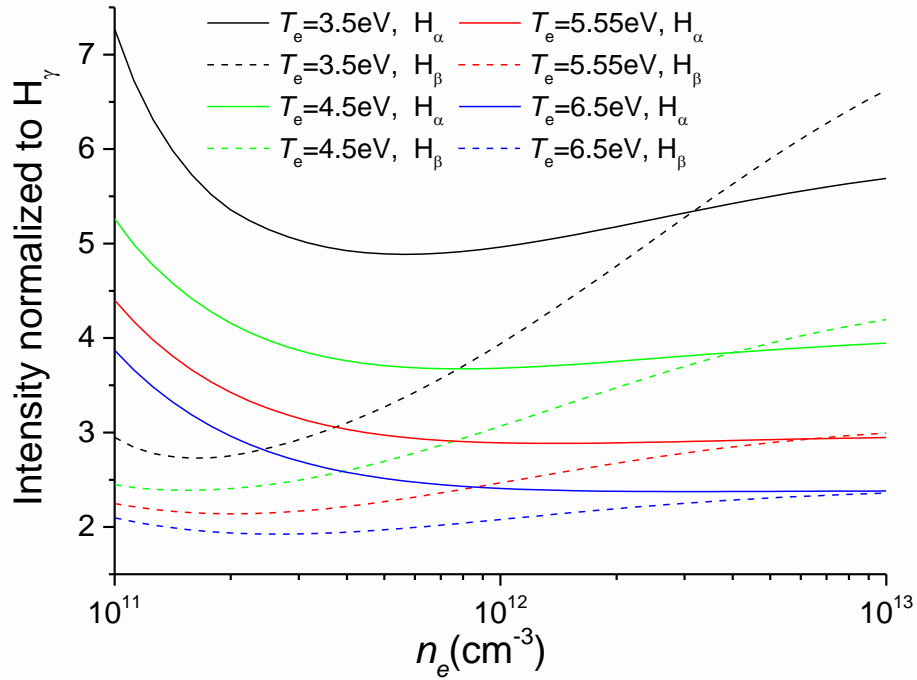
Comparisons of used atomic data



Resulting populations and spectra



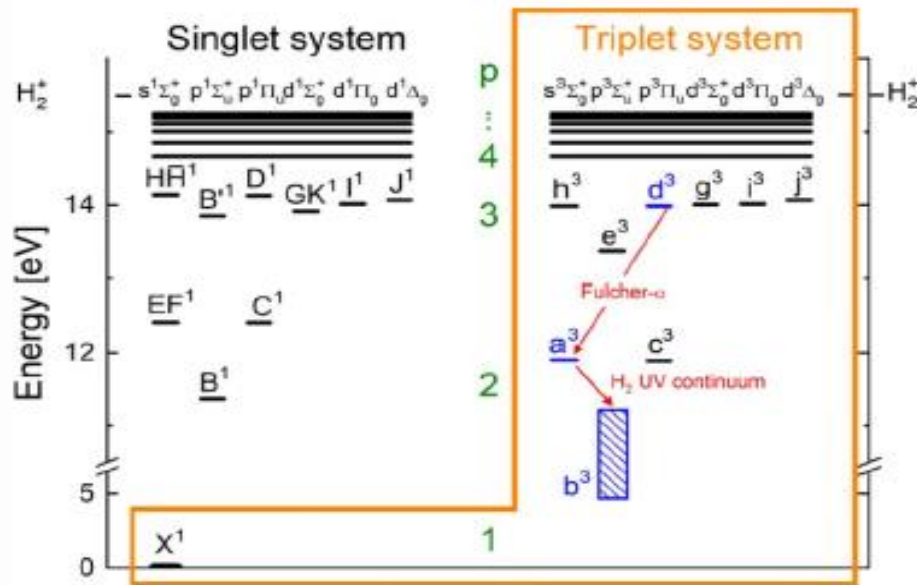
Sensitivity to plasma parameters n_e and T_e



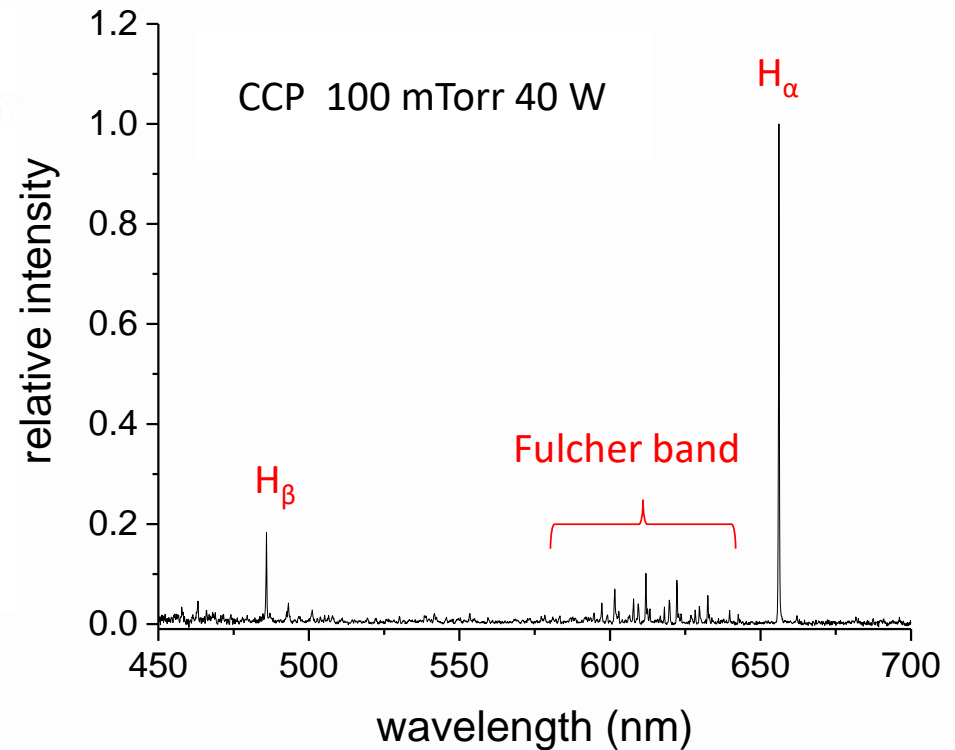
CRM for vibrational states of H molecules



Energy level diagram of H₂



CRM for molecular spectra



Ref. D. Wunderlich et al., J. Phys. D: Appl. Phys. **54** 115201 (2016)

Contents



- Collisional-radiative modeling (CRM) for He, H/D plasmas
 - Modeling methods and AM data
 - Sensitivities to used AM data and plasma parameters
- S/XB ratio for sputtering yield of W
 - S/XB ratio methods and atomic data
 - Experimental S/XB ratios

S/XB ratio for sputtering yield



$$\Gamma = \int n_e S_{Z \rightarrow Z+1} N_Z dx = \int n_e \frac{S_{Z \rightarrow Z+1}}{\frac{n_i}{N_Z} A_{ij}} \left(\frac{n_i}{N_Z} A_{ij} \right) N_Z dx$$

$$= \int \frac{\sum_{\sigma} n_{\sigma} S_{\sigma} / N_Z}{\sum_{\sigma} n_{\sigma} q_{\sigma i} / N_Z \cdot A_{ij} / \sum_{k < i} A_{ik}} n_i A_{ij} dx = \frac{S}{XB} I$$

$$S_{Z \rightarrow Z+1} N_Z = \sum_{\sigma} S_{\sigma} n_{\sigma}$$

$$n_i \sum_{k < i} A_{ik} = n_e \sum_{\sigma} n_{\sigma} q_{\sigma i}$$

$$N_Z = \sum_{\sigma} n_{\sigma},$$

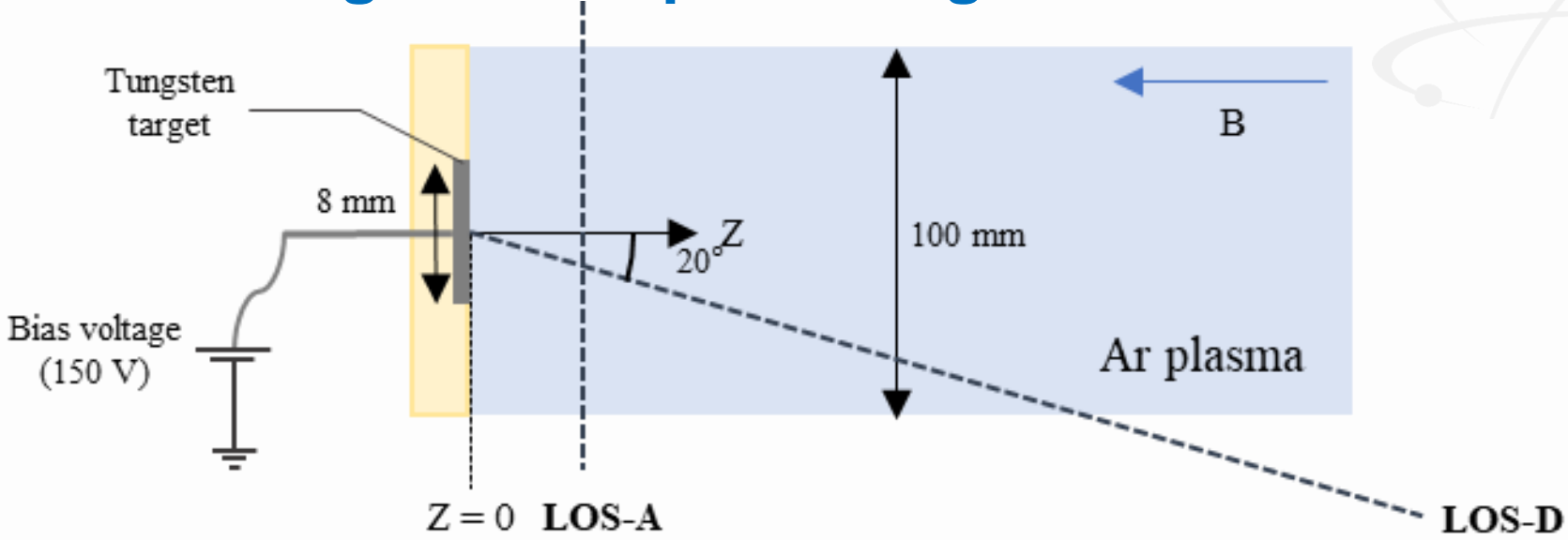
$$\frac{S}{XB} I = \int n_i A_{ij} dx$$

$$\Gamma_W^{\text{Spt}} = \Gamma_{W^+} + \Gamma_W^{\text{GL}}$$

$$\Gamma_W^{\text{Spt}} \approx \Gamma_{W^+} = 4\pi \frac{S}{XB} I_{WI}$$

$$\frac{S}{XB} = \frac{\Gamma_W^{\text{Spt}}}{4\pi I_{WI}} = \frac{Y \Gamma_i}{4\pi I_{WI}}$$

Schematic diagram for optical diagnostic in KPBIF



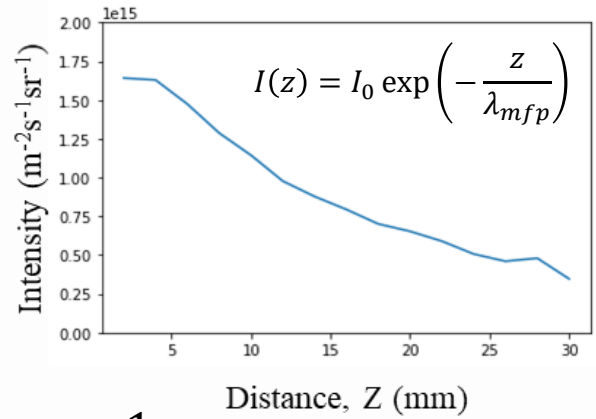
$$\frac{\Gamma_W^{GL}}{\Gamma_W} = \exp\left(-\frac{L}{\lambda_{mfp}}\right)$$

$$\frac{\Gamma_{W^+}}{\Gamma_W} = 1 - \exp\left(-\frac{L}{\lambda_{mfp}}\right)$$

D. Nishijima et al.,
 "Experimental determination
 of S/XB values of W I visible
 lines." Phys. Plasmas **16**
 122503 (2009)

$$\frac{S}{XB} = \frac{\Gamma_{W^+}}{4\pi I_{WI}} = \frac{Y\Gamma_i}{4\pi I_{WI}} \times \left\{ 1 - \exp\left(-\frac{L}{\lambda_{mfp}}\right) \right\}$$

Correction factor for the geometrical loss

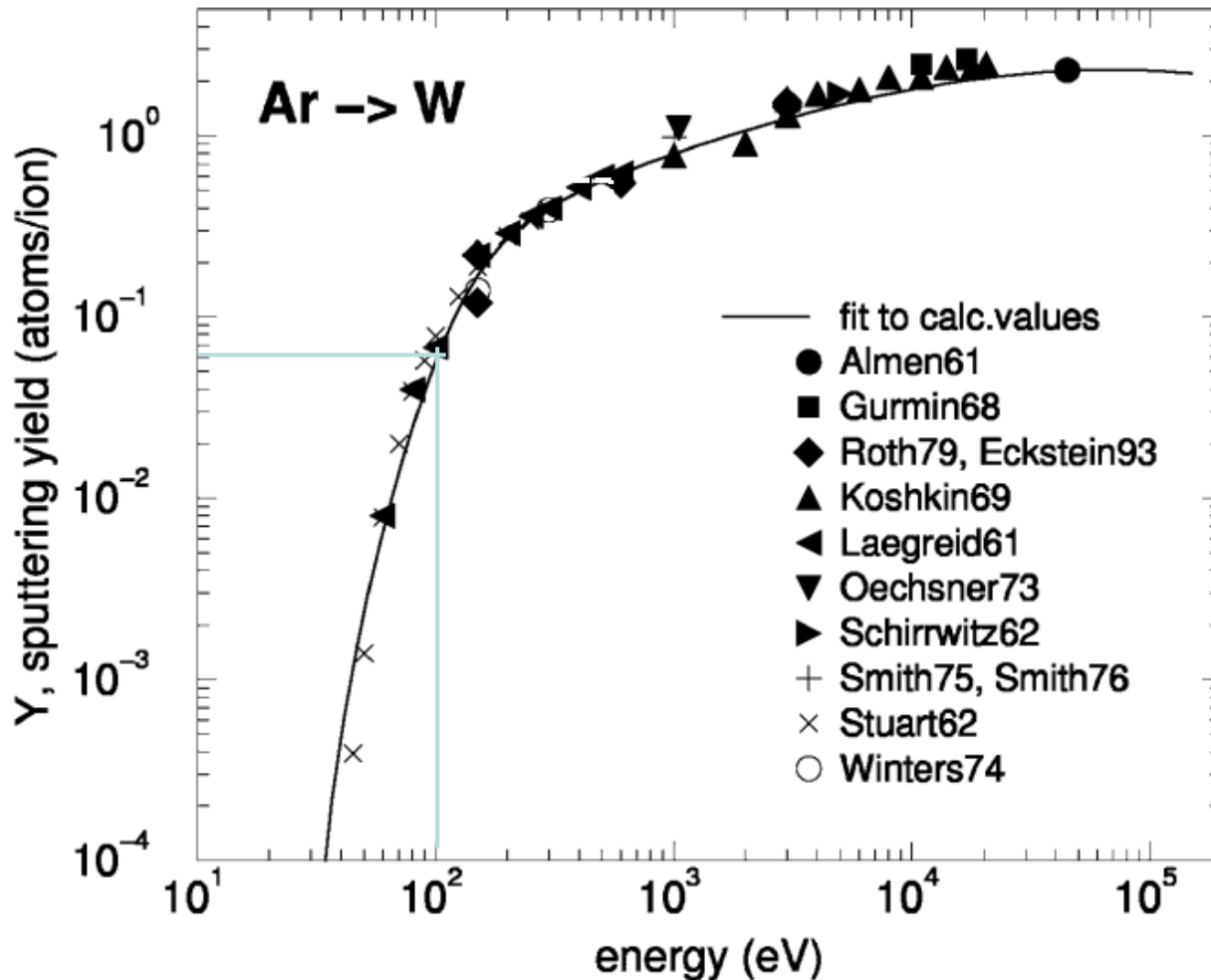


$$\frac{dN_W(x)}{dx} = -\frac{1}{\lambda_{mfp}} N_W(x)$$

$$\frac{dN_W(x)}{dt} = -RC(T_e)n_e N_W(x), \quad \frac{dt}{dx} = \frac{1}{V_W}$$

$$\lambda_{mfp} = \frac{V_W}{RC(T_e)n_e}$$

Sputtering yield of W by Ar ion beam

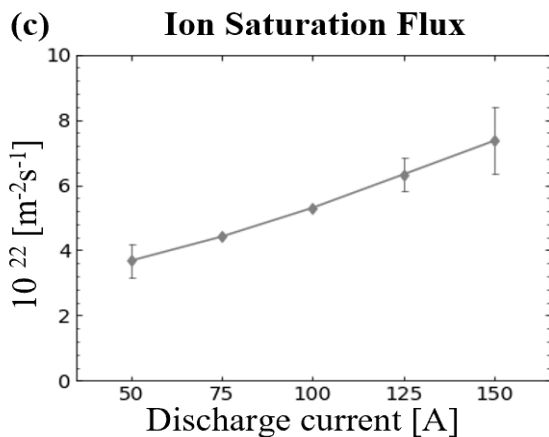
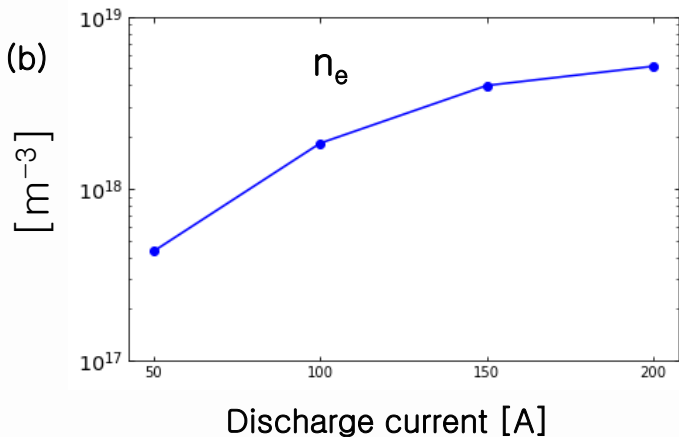
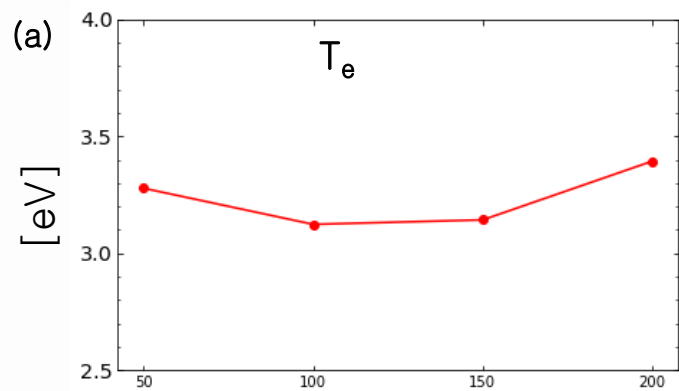


$Y=0.06$ at 100 eV

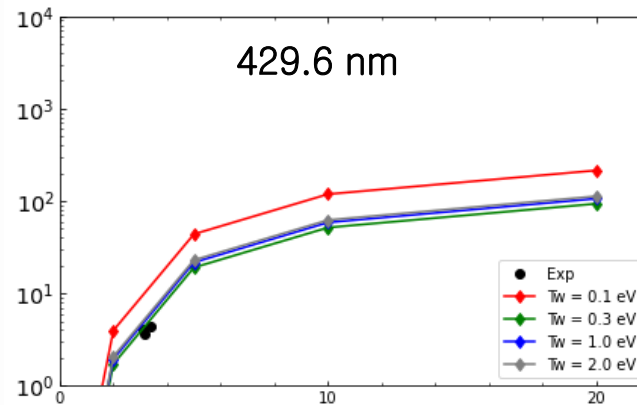
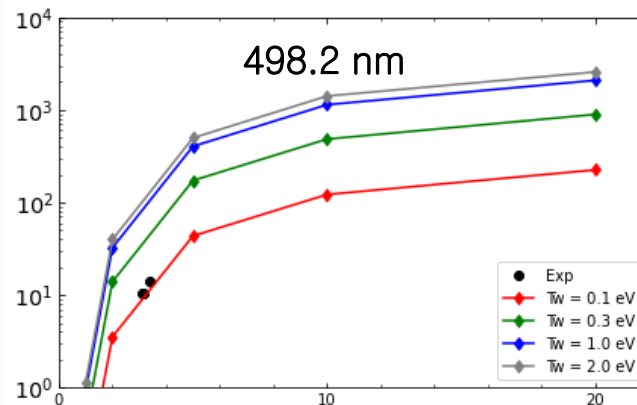
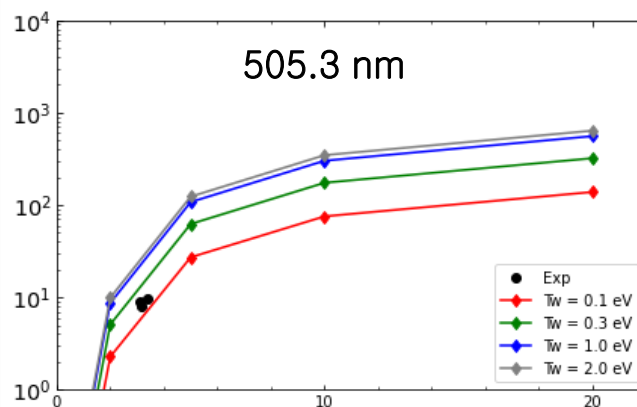
W. Eckstein,
“Sputtering by Particle
Bombardment”,
Springer-Verlag, Berlin
(2007)

Measured S/XB ratios

LP measurements



S/XB ratios : measured vs. calculated



Calculated data

I. Beigman et al., "Tungsten spectroscopy for the measurement of W-fluxes from plasma facing components", Plasma Physics and Controlled Fusion **49** 1833 (2007)

Calculations of S/XB by I. Beigman et al. (2007)

Electron impact ionization (EII) rate

$$\langle v \sigma_{iz} \rangle = 10^{-8} A \frac{\sqrt{\beta}(\beta + 1 + D)}{(\beta + \chi)(\beta + 1)\sqrt{\beta_{iz} + 1}} e^{-\beta_{iz}} (\text{cm}^3 \text{s}^{-1})$$

$$\beta = Ry/T_e, \beta_{iz} = E_{iz}/T_e, E_{iz} = 7.9 \text{ eV}, \quad A = 84.9, \chi = 0.22, D = -0.4$$

for the 6s and 5d shells by Born-Ochkur approximation using ATOM

Electron impact excitation (EIE) rate

$$\langle v \sigma_{k_0 k} \rangle = 0.11 \times 10^{-16} \frac{g_k}{g_{k_0}} A_{kk_0} \left(\frac{Ry}{\Delta E} \right)^3 u(T_e) e^{-\beta_{ex}} (\text{cm}^3 \text{s}^{-1})$$

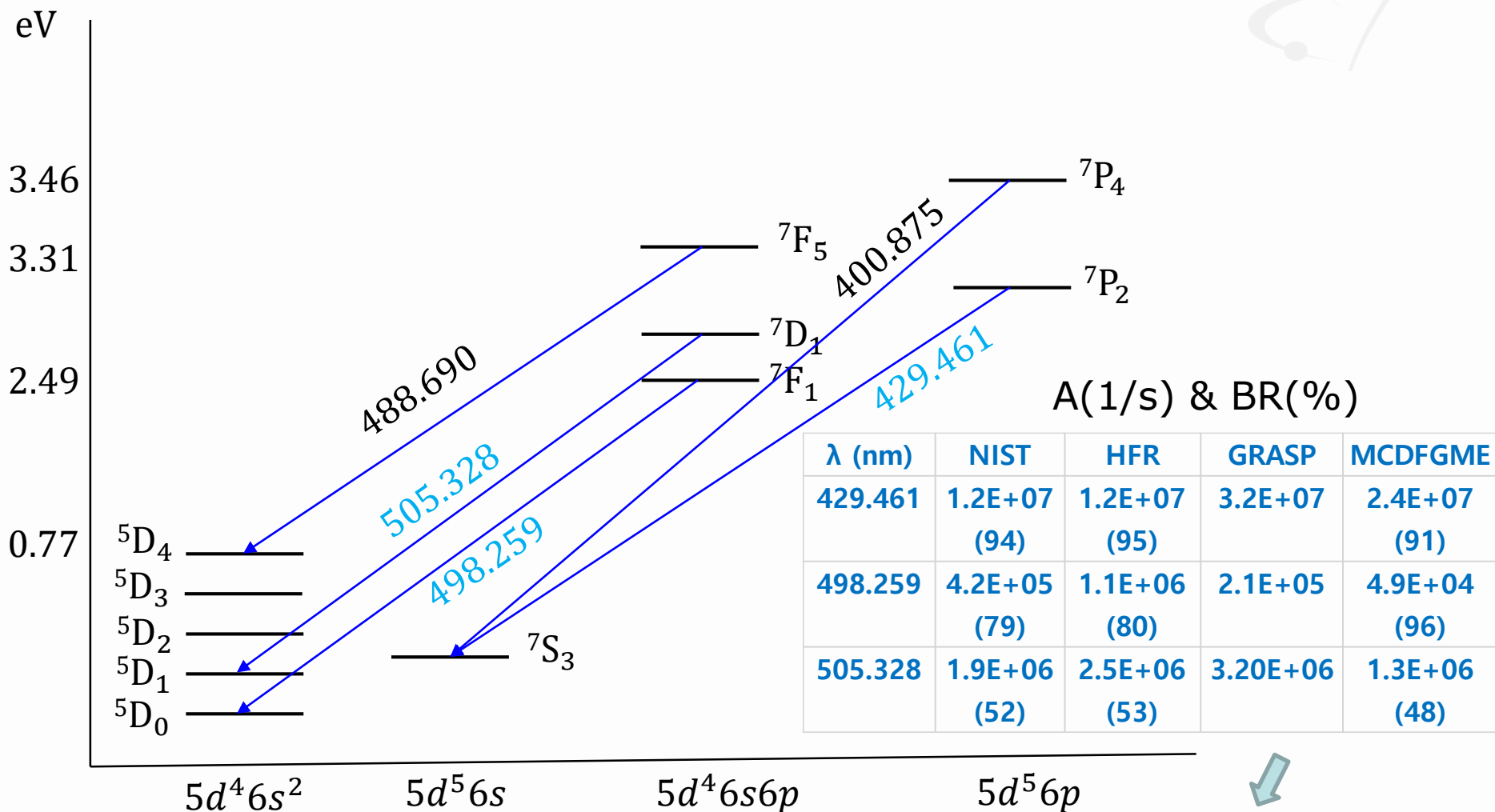
$$u(T_e) = \beta^{0.5} \log \left(2 + \frac{1}{1.78\beta_{ex}} \right), \quad \beta = Ry/T_e, \quad \beta_{ex} = \Delta E/T_e$$

for the dipole transitions by the semi-empirical van Regemorter formula (Astrophys. J. **136** 906 (1962))

The ionization per photon

$$\frac{S}{XB} = \langle v \sigma_{iz} \rangle / \sum_{k_0} N_{k_0} \langle v \sigma_{k_0 k} \rangle \cdot \sum_{k''} A_{kk''} / A_{kk'}$$

Selected lines for S/XB of W I



NIST : A. E. Kramida and T. Shirai , J. Phys. Chem. Ref. Data **35** 423 (2006)

HFR : P. Quinet et al., J. Phys. B: At. Mol. Opt. Phys. **44** 145005 (2011), DESIRE DB

GRASP : R. T. Smyth, C. P. Ballance et al., Phys. Rev. A **97** 052705 (2018)

MCDFGME : Present calculation

Energy levels for S/XB of W I



Energies (eV)

Configuration	Term	J	NIST	HFR	GRASP	MCDFGME
5d ⁴ 6s ²	⁵ D	0	0	0	0	0
		1	0.2071	0.2104	0.1273	0.1465
		2	0.4123	0.4192	0.2927	0.3239
		3	0.5988	0.5998	0.4729	0.5204
		4	0.7711	0.7585	0.6631	0.6941
5d ⁵ 6s	⁷ S	3	0.3659	0.3587	0.4057	0.0751
5d ⁴ 6s6p	⁷ F	1	2.4877	2.4826	2.1842	2.1667
	⁷ D	1	2.6599	2.6581	2.4292	2.4960
5d ⁵ 6p (5d ⁴ 6s6p)	⁷ P	2	3.2521	3.2584	2.9697	2.9706

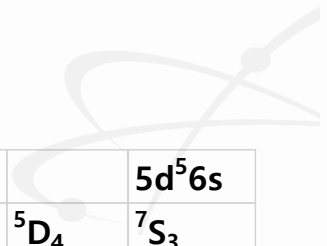
Configurations for energies W I

NIST, HF based on EXP (Kramida et al. 2006)	HFR (Wyart 2010)	Even parity		
		HFR (Quinet et al. 2011)	GRASP (Smyth et al. 2018)	MCDFGME (Present)
$5d^46s^2$	$5d^46s^2$	$5d^46s^2$	$5d^46s^2$	$5d^46s^2$
$5d^56s$	$5d^56s$	$5d^56s$	$5d^56s$	$5d^56s$
$5d^6$	$5d^6$	$5d^6$	$5d^6$	$5d^6$
$5d^46p^2$		$5d^46p^2$	$5d^46p^2$	$5d^46p^2$
		$5d^46d^2$	$5d^46d^2$	$5d^46d^2$
$5d^36s6p^2$		$5d^36s6p^2$	$5d^36s6p^2$	$5d^36s6p^2$
				$5d^36s6d^2$
		$5d^26s^26p^2$	$5d^26s^26p^2$	$5d^26s^26p^2$
			$5d^26s^26d^2$	$5d^26s^26d^2$
			$5d^47s^2$	$5d^47s^2$
$5d^46s6d$		$5d^46s6d$	$5d^46s6d$	
		$5d^56d$	$5d^56d$	
$5d^46s7s$		$5d^46s7s$	$5d^46s7s$	
$5d^57s$		$5d^57s$	$5d^57s$	$5d^57s$
$5d^58s$				
$5d^46s7d$				
			$5d^36s^26d$	
			$5d^36d^3$	
			$5d^36s^27s$	
			$5d6s^26d^3$	
			$5p^45d^8$	
			$5p^45d^66s^2$	
			$5p^45d^76s$	

NIST, HF based on EXP (Kramida et al. 2006)	HFR (Wyart 2010)	Odd parity		
		HFR (Quinet et al. 2011)	GRASP (Smyth et al. 2018)	MCDFGME (Present)
$5d^46s6p$	$5d^46s6p$	$5d^46s6p$	$5d^46s6p$	$5d^46s6p$
$5d^56p$	$5d^56p$	$5d^56p$	$5d^56p$	$5d^56p$
$5d^36s^26p$	$5d^36s^26p$	$5d^36s^26p$	$5d^36s^26p$	$5d^36s^26p$
		$5d^36p^3$	$5d^36p^3$	$5d^36p^3$
		$5d^26s6p^3$	$5d^26s6p^3$	$5d^26s6p^3$
		$5d^46s7p$		
		$5d^57p$		
		$5d^46s5f$		
		$5d^55f$		
			$5d^36s^27p$	
			$5p^55d^7$	
			$5p^55d^66s$	

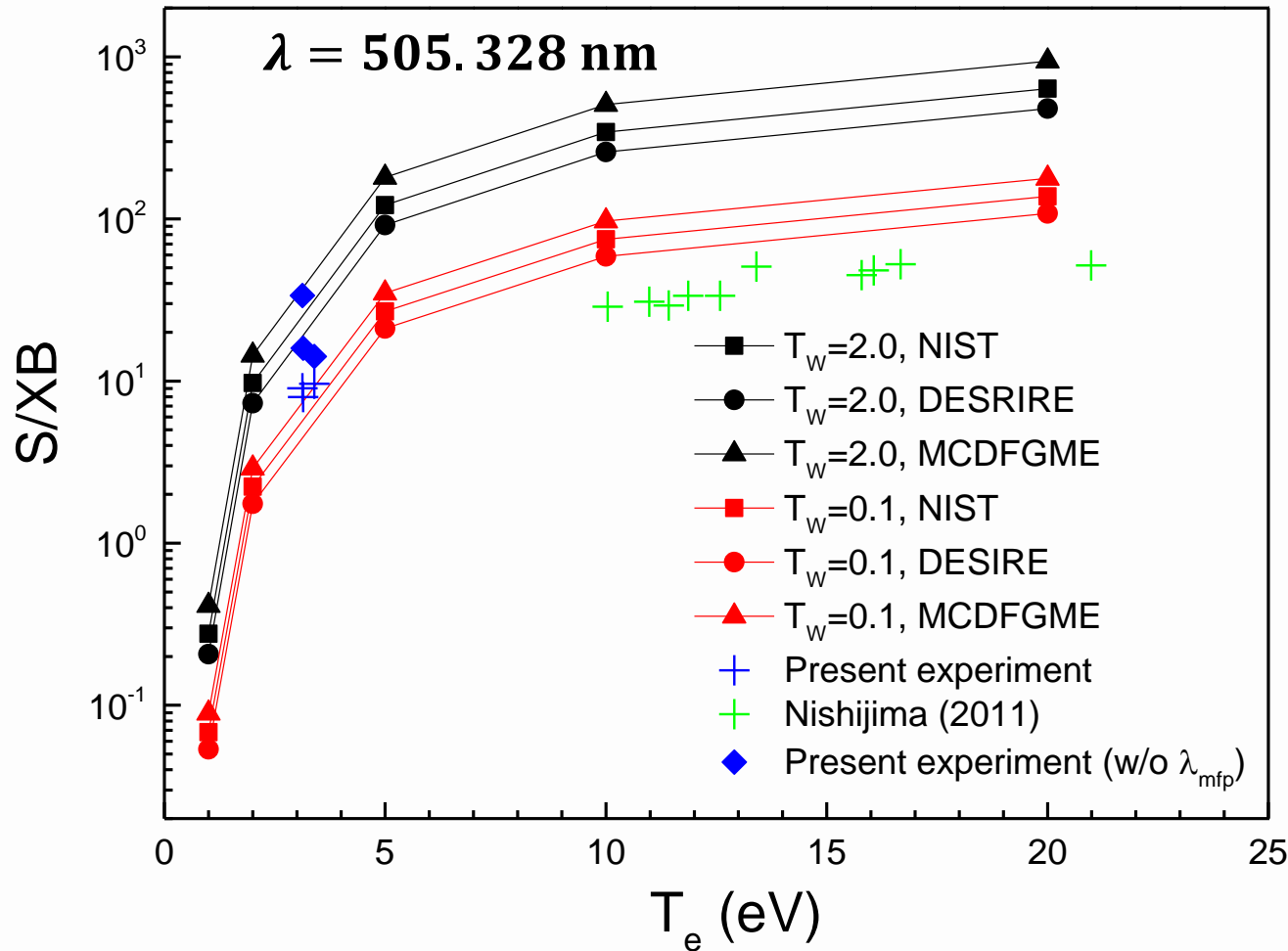
- ❖ Our present MCDFGME : **Self-consistent-field (SCF) radial wavefunctions in the configuration interaction (CI) procedure**
- ❖ Others : **Fixed radial wavefunctions in the CI procedure, (GRASP : Extended Average Level (EAL) method)**

S/XB for various DB's



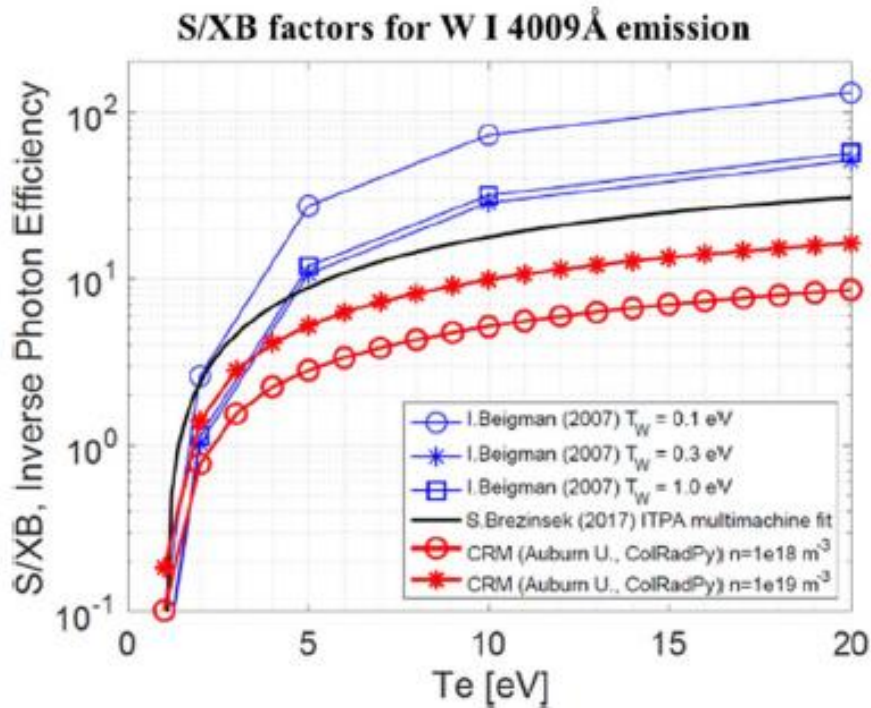
$$N_{k_0} = \frac{(2J_{k_0} + 1) \exp(-\Delta E_{k_0}/T_W)}{\sum_k (2J_k + 1) \exp(-\Delta E_k/T_W)}$$

Config.	5d ⁴ 6s ²					5d ⁵ 6s
Level	⁵ D ₀	⁵ D ₁	⁵ D ₂	⁵ D ₃	⁵ D ₄	⁷ S ₃
N _{k₀} , T _w =2eV	0.04	0.11	0.16	0.21	0.25	0.23
N _{k₀} , T _w =0.1eV	0.60	0.23	0.05	0.01	0.00	0.11

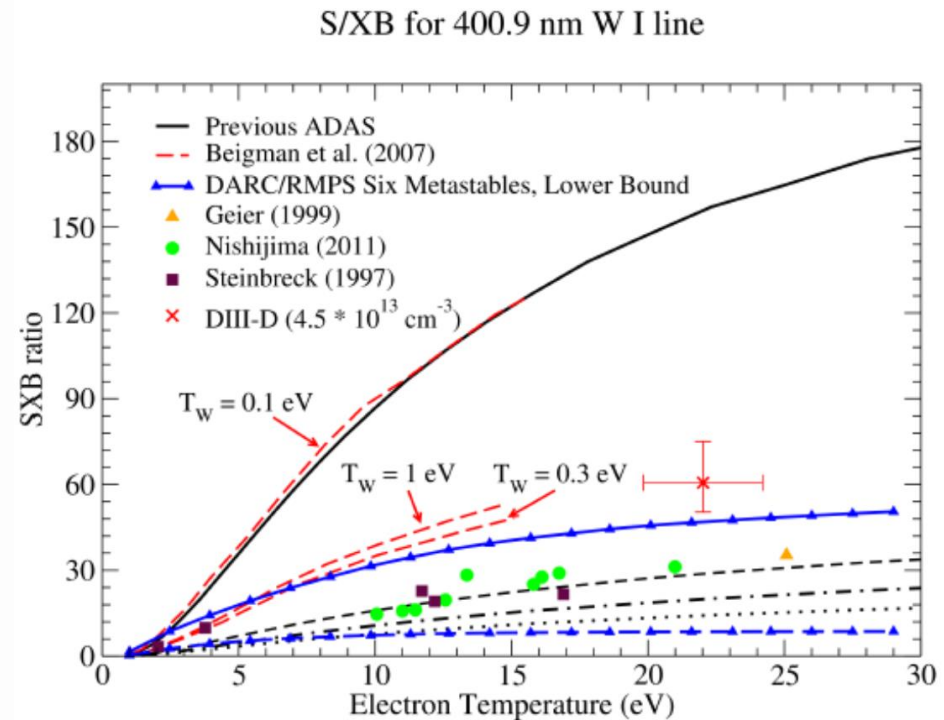


D. Nishijima et al., *Phys. Plasmas* **18** 019901 (2011),
 Erratum of *Phys. Plasmas* **16** 122503 (2009)

S/XB by EIE cross section from DARC calculation and a CRM



C. C. Kepler et al.,
 Plasma Phys. Control.
 Fusion **64** 104008
 (2022)



C. J. Favreau,
 Ph. D thesis,
 Auburn Univ. (2019)

Further effects to be considered



Opacity, Optical trapping

$$\frac{S}{XB} = \frac{\eta \Gamma_{W+}}{4\pi I_{WI}} = \frac{\eta Y \Gamma_i}{4\pi I_{WI}} \times \left\{ 1 - \exp\left(-\frac{L}{\lambda_{mfp}}\right) \right\}, \quad \eta < 1.0$$

Optical cascade from upper levels

Finally, more accurate collisional-radiative modeling considering other processes such as diffusion etc.



Other models for EII(E) cross sections

Binary Encounter Bethe (BEB) model for EII

(Y.-K. Kim and M. E. Rudd, Phys. Rev. A **50** 3954 1994)

$$\sigma_{\text{neut}} = \frac{4\pi a_0^2 N R^2}{B^2 [t + (u + 1)/m]} \left[\frac{\ln t}{2} \left(1 - \frac{1}{t^2} \right) + 1 - \frac{1}{t} - \frac{\ln t}{t + 1} \right]$$

N : The orbital occupation number

B : The orbital binding energy,

$t = T/B$ for incident electron energy T

$u = U/B$ for the orbital kinetic energy U

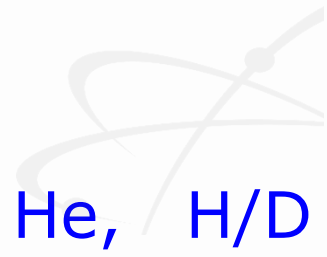
Constant $m = 1$ for K- and L-shell, $m=n$ of other orbitals

BE scaled plane wave Born (PWB) cross section for EIE

(Y.-K. Kim, Phys. Rev. A **64** 032713 2001)

$$\sigma_{\text{BE}} = \frac{T}{T + B + E} \sigma_{\text{PWB}} \quad E : \text{The excitation energy}$$

Summary and outlook



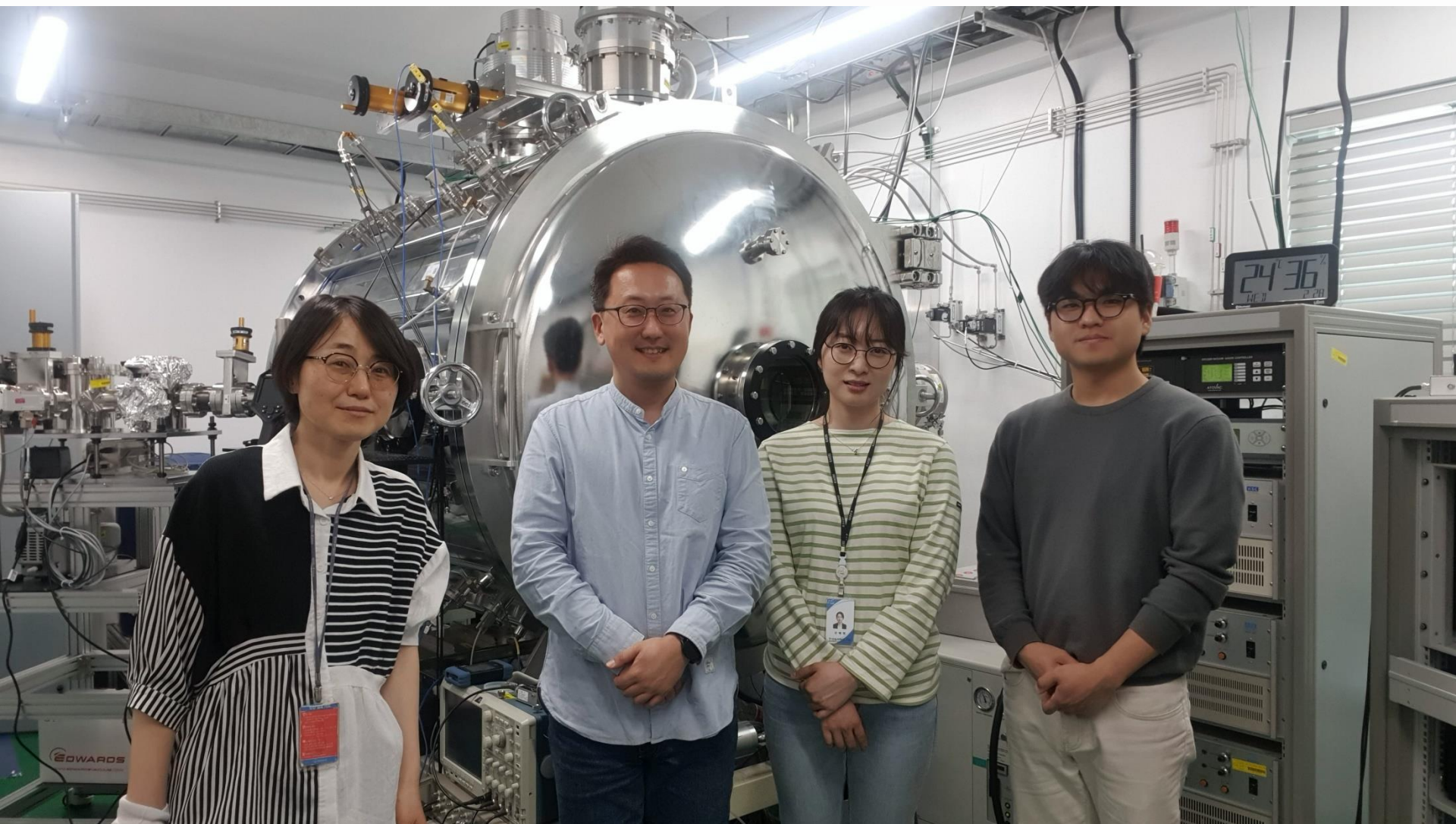
- Collisional-radiative modeling (CRM) for He, H/D plasma
- ✓ The CRM solves nonlinear steady state balance equations including processes such as radiation trapping and heavy particle collisions self-consistently for low temperature regime. Sensitivities of line spectra and particle densities to used AM data and plasma parameters are investigated.
- ✓ Charge exchanges such as $H^+ + H_2 \rightarrow H + H_2^+$ etc. will be considered in the CRM for H.
- ✓ CRM for H taking into account for molecular vibrational states will be performed.

Summary and outlook



- S/XB ratio for sputtering yield of W
- ✓ S/XB ratios for W I were spectroscopically measured in KPBI-F and compared with model calculations using EII/EIE and radiative transition data.
- ✓ New atomic data by metastable resolved BEB cross section for EII and BE scaled PWB cross section for EIE combined with MCDF calculation will be used for the S/XB calculation.

Our group members



Acknowledgements



- **Wonho Choe**, Gas discharge physics lab., Korea Advanced Institute of Science and Technology, Daejeon, Rep. of Korea
- **Octavio Roncero**, Quantum molecular dynamics, Instituto de Física Fundamental, Madrid, Spain
- **Dmitry Fursa**, MCCC molecular data, Curtin Institute for Computation and Department of Physics, Astronomy and Medical Radiation Sciences, Curtin University, Australia